

# Changes in Stratospheric Climate and Age-of-Air in Recent GEOS Systems since MERRA-2

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## Key Points:

- The stratospheric mean age-of-air simulated in GEOS is sensitive to the remapping scheme used within the finite-volume dynamical core.
- This sensitivity in the age-of-air is significant ( $\sim 30\%$ ) and imprints on the simulated distributions of several long-lived chemical trace gases, including nitrous oxide and methane.
- The age-of-air sensitivities primarily reflect changes in resolved wave convergence over the Northern Hemisphere midlatitude stratosphere, which impact mean upwelling in the tropical lower stratosphere.

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18 **Abstract**

19 Accurately modeling the large-scale transport of trace gases and aerosols is critical for interpreting past (and projecting future) changes in atmospheric composition. 20 Simulations of the stratospheric mean age-of-air continue to show persistent biases in chemistry climate models, although the drivers of these biases are not well understood. Here 21 we identify one driver of simulated stratospheric transport differences among various NASA 22 Global Earth Observing System (GEOS) candidate model versions under consideration 23 for the upcoming GEOS Retrospective analysis for the 21<sup>st</sup> Century (GEOS-R21C). In 24 particular, we show that the simulated age-of-air values are sensitive to the so-called “remap- 25 ping” algorithm used within the finite-volume dynamical core, which controls how in- 26 dividual material surfaces are vertically interpolated back to standard pressure levels af- 27 ter each horizontal advection time step. Differences in the age-of-air resulting from changes 28 within the remapping algorithm approach  $\sim 1$  year over the high latitude middle strato- 29 sphere – or about 30% climatological mean values – and imprint on several trace gases, 30 including methane ( $\text{CH}_4$ ) and nitrous oxide ( $\text{N}_2\text{O}$ ). These transport sensitivities reflect, 31 to first order, changes in the strength of tropical upwelling in the lower stratosphere (70- 32 100 hPa) which are driven by changes in resolved wave convergence over northern mid- 33 latitudes as (critical lines of) wave propagation shift in latitude. Our results strongly sup- 34 port continued examination of the role of numerics in contributing to transport biases 35 in composition modeling. 36 37

38 **Plain Language Summary**

39 Large-scale transport plays a crucial role in distributing climatically important trace 40 constituents in the atmosphere, especially in the stratosphere where transport largely 41 determines the chemical lifetimes of trace gases. One summary of transport in the strato- 42 sphere is the “mean age” or the mean transit time since air at a point in the stratosphere 43 was last in the troposphere. Current models used for simulating stratospheric compo- 44 sition produce a range of simulated ages, although these differences are poorly under- 45 stood. Among other factors, model numerics play a critical role in transport, but few 46 studies have explored the sensitivity of the mean age to the choice of numerical scheme 47 employed within different dynamical cores. Here we use one model to show that the mean 48 age is sensitive to the so-called “remapping” algorithm used within the finite-volume dy- 49 namical core that controls how individual material surfaces are vertically interpolated 50 back to standard pressure levels after each horizontal advection time step. This reflects 51 sensitivities in the representation of how waves propagate from the troposphere into the 52 stratosphere. This work suggests that model numerics can be an important factor in con- 53 tributing to differences in simulated transport among models.

54 **1 Introduction**

55 The chemical and radiative properties of the troposphere and lower stratosphere 56 are strongly influenced by the stratosphere-troposphere exchange of mass and tracers (e.g., 57 Morgenstern and Carver (2001); Hegglin et al. (2006); Pan et al. (2007)). Properly sim- 58 ulating the stratospheric circulation and its influence on atmospheric composition in earth 59 system models is important for capturing past decadal trends in surface climate, par- 60 ticularly in response to changes in Southern Hemisphere ozone depletion (e.g., Son et 61 al. (2009); Polvani et al. (2011)). In the Northern Hemisphere (NH), the stratospheric 62 circulation’s coupling to ozone could represent an important feedback on the climate’s 63 response to future increases in greenhouse gases (GHGs), especially over the North At- 64 lantic (e.g., Chiodo and Polvani (2019)). On shorter subseasonal timescales, stratospheric 65 ozone changes associated with strong polar vortex states may also modulate Arctic sea 66 level pressure and surface temperatures (e.g., Ivy et al. (2017); Oehrlein et al. (2020)), 67 so much so that seasonal forecast systems employing prognostic ozone show suggestions

68 of increased signal-to-noise ratio in predictions of the North Atlantic Oscillation (B. M. Monge-  
69 Sanz et al. (2022)).

70 Key to accurately simulating a consistent representation of coupling between strato-  
71 spheric dynamics and chemical trace gases is ensuring that a model’s underlying trans-  
72 port circulation is properly represented. To this end, much effort has been paid to de-  
73 veloping and refining so-called “tracer-independent” metrics of transport (Holzer and Hall  
74 (2000)) such as the mean age-of-air (Hall and Plumb (1994)) and to applying these mea-  
75 sures to rigorously evaluate model transport characteristics in chemistry climate mod-  
76 els (CCMs) (e.g., Hall et al. (1999); Orbe et al. (2018); Dietmüller et al. (2018); Aba-  
77 los et al. (2020)).

78 While the assessment of CCMs participating in the SPARC Chemistry Climate Model  
79 Validation (SPARC CCMVal) effort showed a marked improvement in simulated trans-  
80 port characteristics relative to previous intercomparisons (J. Neu et al. (2010)), more re-  
81 cent analysis of models participating in the SPARC Chemistry Climate Modeling Ini-  
82 tiative (CCMI) (Eyring et al. (2013)) do not demonstrate any improvement (Dietmüller  
83 et al. (2018), see their Figure 3). In particular, although some models produce mean age  
84 values that agree well with observational estimates, the CCMI intermodel spread is  $\sim$   
85 50%, with models generally simulating transport that is too vigorous relative to obser-  
86 vations. While documenting these transport differences among models is straightforward,  
87 understanding the drivers of this spread remains a challenge and there is still no con-  
88 sensus on what is causing the large spread in simulated ages among the current gener-  
89 ation of CCMs.

90 A key challenge in identifying the drivers of age-of-air – and other stratospheric trans-  
91 port – biases is that they reflect the time-integrated effects of advection by the residual  
92 mean circulation and eddy diffusive mixing, or the quasi-random transport due to the  
93 breaking of Rossby waves (e.g., Holton et al. (1995); Plumb (2002)). Given that the in-  
94 fluences of mixing and advection are not easily separable, studies have come to differ-  
95 ent conclusions about sources of age biases in models. In particular, the analysis of the  
96 CCMVal models showed a strong correlation between the intermodel spread in the age-  
97 of-air and lower stratospheric tropical upwelling, whereas Dietmüller et al. (2018) showed  
98 that the age spread among the CCMI models was driven by differences in mixing. While  
99 future attempts to further distinguish between sources of age biases using either simpli-  
100 fied “leaky pipe” models (Plumb (1996); J. L. Neu and Plumb (1999)) or more complete  
101 measures of the transport circulation such as the “age spectrum” (e.g., Hall and Plumb  
102 (1994); Waugh and Hall (2002))) may prove enlightening, at present there is no consen-  
103 sus on what is causing large simulated age-of-air biases in models.

104 One potential limitation of previous work based on multi-model intercomparisons  
105 is that many aspects of model formulation can influence both stratospheric upwelling and  
106 mixing. Thus, while intercomparisons are useful for identifying common model biases,  
107 understanding the drivers of these biases is difficult absent single model-based process  
108 studies. Among these, several aspects of model formulation have been identified as in-  
109 fluencing simulated mean age distributions. As the mean age is sensitive to vertical mo-  
110 tion in the lowermost stratosphere, these include large sensitivities to vertical resolution  
111 (Orbe et al. (2020)) and to spurious vertical mixing either introduced in vertical coor-  
112 dinate transformations in offline chemical transport models (B. Monge-Sanz et al. (2007))  
113 or through use of assimilated winds performed either in offline (e.g., Legras et al. (2004))  
114 or online data assimilation and “nudged” configurations (e.g., Pawson et al. (2007); Orbe  
115 et al. (2017); Davis et al. (2022)). These age sensitivities can be still further amplified,  
116 depending on whether or not parameterized gravity waves are included (Eichinger et al.  
117 (2020)).

118 By comparison, sensitivities of the mean age to underlying tracer numerics have  
119 been less well examined, although Eluszkiewicz et al. (2000) documented a large sensi-

120 tivity in simulated age-of-air values to the choice of advection scheme. More recently,  
 121 Gupta et al. (2020) showed differences of  $\sim 25\%$  in the age-of-air across identical exper-  
 122 iments performed using four different dynamical cores, especially between those using  
 123 spectral versus finite-volume schemes. The experiments employed in that study, how-  
 124 ever, were highly idealized and it is not clear if the strong influence of tracer numerics  
 125 that they identified is also realized in more comprehensive model simulations with moist  
 126 physics, especially in the context of model development as carried out in operational mod-  
 127 eling centers.

128 To better elucidate this influence of tracer numerics on the transport properties sim-  
 129 ulated in a comprehensive global model context, here we document the sensitivity of the  
 130 stratospheric mean age in several recent versions of the NASA Global Earth Observing  
 131 System (GEOS) general circulation model (Molod et al., 2015). The versions examined  
 132 here represent different stages in model development since the Modern-Era Retrospec-  
 133 tive Analysis for Research and Applications Version 2 (MERRA-2; Gelaro et al. (2017)).

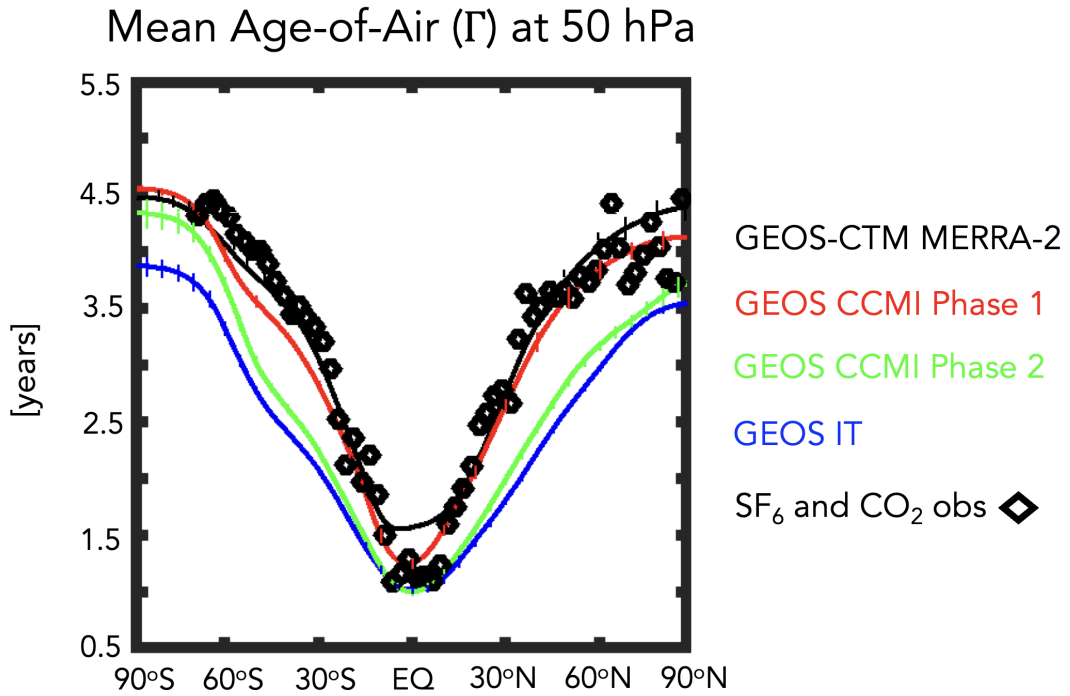
134 Our focus on transport evaluation is partly in wake of the upcoming release of the  
 135 GEOS Retrospective analysis for the early 21<sup>st</sup> Century (GEOS-R21C), which is an at-  
 136 mospheric reanalysis that includes many advances over MERRA-2, and serves as a step  
 137 towards MERRA-3, a planned coupled Earth system reanalysis. Along with having im-  
 138 plications for GEOS-R21C, our analysis also includes the GEOS-IT Version 5.29.4 dat-  
 139 ing from October 19, 2022 (hereafter simply “GEOS-IT”), which will be used to drive  
 140 an off-line chemistry reanalysis with a full chemistry model (GEOS-Chem) and an ad-  
 141 vanced Constituent Data Assimilation component to update the chemistry fields. Since  
 142 this chemical reanalysis will be produced in replay-mode (one-way coupling) whereby  
 143 the meteorology fields are used to define the background atmospheric flow (Orbe et al.,  
 144 2017), it is imperative that recent tags of GEOS produce a credible representation of trans-  
 145 port processes.

146 In particular, here we document how in the process of evaluating candidate sys-  
 147 tems for GEOS-R21C we found that the mean age was  $\sim 1$  year younger (or  $\sim 30\%$  smaller)  
 148 than the values simulated in the model version used to produce MERRA-2 (Figure 1).  
 149 The model versions shown in Figure 1 reflect more than 10 years’ worth of accumulated  
 150 changes in model development, most notably changes in radiation, parameterized con-  
 151 vection and, as we focus on here, changes in the model’s description of vertical advec-  
 152 tion. More precisely, we show that slight modifications to the so-called “remapping” al-  
 153 gorithm, which is used to transform advected fields from Lagrangian levels to the new  
 154 pressure levels after each horizontal advection time step, are the primary driver of the  
 155 age-of-air changes exhibited in recent GEOS-R21C candidate model versions. Thus, whereas  
 156 Gupta et al. (2020) highlighted large differences between dynamical cores employing spec-  
 157 tral versus finite-volume (FV) numerics, our results show that large transport differences  
 158 can occur even within a given FV dynamical core, a result which may have broader im-  
 159 plications for other general circulation models employing FV numerics. We begin by dis-  
 160 cussing methods in Section 2 and present key results and conclusions in Sections 3 and  
 161 4, respectively.

## 162 2 Methods

### 163 2.1 Model Configurations

164 Here we present results from several versions of GEOS spanning MERRA-2 to more  
 165 recent candidates for GEOS-R21C. Among these model versions, a subset are more “of-  
 166 ficial” as they have been documented and/or employed in recent model intercomparisons  
 167 and are highlighted in Figure 1. In particular, these include a model version that was  
 168 used in Phase 1 of CCM1 and documented in Orbe et al. (2017) (Fig. 1, red line). A more



**Figure 1.** The 2000-2010 climatological annual mean meridional profile of the stratospheric mean age-of-air ( $\Gamma$ ), evaluated at 50 hPa. Results from a GEOS-CTM integration constrained with MERRA-2 meteorological fields (black line) as well as free-running GEOS simulations using a model configuration for CCM1 Phase 1 (red line), CCM1 Phase 2 (green line) and the GEOS-IT Version 5.29.4 dating from October 19, 2022 (hereafter simply "GEOS-IT", blue line) are shown. All simulations are constrained with the same (observed) historical sea surface temperatures. Diamonds correspond to SF<sub>6</sub> and CO<sub>2</sub> in situ based estimates of  $\Gamma$  from Boering et al. (1996) and Engel et al. (2009). Vertical dashed lines denote  $\pm\sigma$ , the standard deviation of  $\Gamma$  over 2000-2010, for each model simulation.

169 recent model version that was used in the CCM1 Phase 2 simulations (correspondence  
170 with Michael Manyin) is also shown (Fig. 1, green line).

171 We begin by comparing 10-year (2000-2010) climatological mean zonally averaged  
172 age-of-air profiles at 50 hPa across this subset of model versions, derived from 30-year  
173 long atmosphere-only (AMIP) integrations constrained with observed sea surface tem-  
174 peratures (Figure 1). First, we note that the profiles for the CCM1 Phase 1 version of  
175 the model are very close to observations (black stars), consistent with the  
176 “GEOSCCM” documented age characteristics reported in Dietmüller et al. (2018) (see  
177 their Figure 3). In addition, while passive tracers were not integrated within MERRA-  
178 2, results using the GEOS chemistry transport model (GEOS-CTM, Kouatchou et al.  
179 (2015)) constrained with MERRA-2 meteorological fields (black line) also exhibits good  
180 agreement with observed values. This good agreement between the CTM-generated age-  
181 of-air and the observations is consistent with results from a previous GEOS-CTM sim-  
182 ulation (constrained with MERRA) as documented in Orbe et al. (2017).

183 Moving to more recent development versions of the model (green and blue lines),  
184 however, reveals a reduction in the mean age by  $\sim 1$  year over both southern and north-  
185 ern high extratropical latitudes, or a decrease of  $\sim 20\text{-}30\%$  relative to the MERRA-2 con-  
186 strained simulation and the observations. As discussed earlier, the green line refers to  
187 the CCM1 Phase 2 model version, whereas the blue line refers to GEOS-IT. Note that  
188 this decrease in the climatological age in both model versions far exceeds the (internal)  
189 variations in mean age that occur interannually (vertical bars on solid lines).

190 Since MERRA-2, numerous updates have been introduced in to the GEOS model.  
191 This includes replacing the Chou Suarez (Chou & Suarez, 1994) and Chou (Chou, 1990,  
192 1992) radiation codes with the Rapid Radiative Transfer Model for GCMs (RRTMG;  
193 Iacono et al. (2008), which impacts the stratosphere; the introduction of the Grell-Freitas  
194 deep-convection code (Grell and Freitas (2014); Freitas et al. (2018)), which has a min-  
195 imal impact on this study; and upgrades to the vertical remapping option from that used  
196 in MERRA-2 (hereafter denoted as GMAO FV3 Cubic) to the current updated GFDL  
197 remapping option (hereafter denoted as GFDL FV3 Cubic). Note that both MERRA-  
198 2 and current systems use the GFDL FV3 dynamical core (Lin et al., 2017), but differ  
199 in terms of this remapping option.

200 As we show in Section 3, the latter transition from the remapping scheme used in  
201 MERRA-2 (GMAO FV3) to the scheme used in the current GFDL FV3 core has the largest  
202 impact on the simulated age-of-air values. This degradation in simulated transport within  
203 the current GFDL FV3 core primarily reflects differences in the implementation of the  
204 remapping algorithm, which is used to vertically interpolate individual material surfaces  
205 after each horizontal advection step back to the model’s reference Eulerian coordinate.  
206 In its implementation in MERRA-2 within the GMAO FV3 core (Table 1), this remap-  
207 ping involves 1) fitting piecewise parabolic (hereafter PPM) functions to input layer-mean  
208 values of  $U$ ,  $V$ ,  $Q$  and tracers; 2) calculating PPM functions to output layer edges; 3)  
209 integrating PPM functions between output layer edges to produce new layer-mean val-  
210 ues of  $U$ ,  $V$ ,  $Q$  and tracers; 4) calculating total energy (TE) at input mid-layer pressures;  
211 5) calculating TE at output mid-layer pressures using cubic interpolation and applying  
212 an a-posteriori integral conservation; and, finally, 6) remapping temperatures from to-  
213 tal energy via  $T = (TE - K - \Phi)/C_p$ . Here  $T$ ,  $U$ ,  $V$ ,  $Q$ ,  $C_p$ ,  $K$  and  $\Phi$  correspond to tem-  
214 perature, zonal wind, meridional wind, specific humidity, specific heat capacity and ki-  
215 netic and potential energy, respectively.

216 Differences in the implementation of this remapping algorithm within the current  
217 GFDL FV3 core results in degraded simulated stratospheric transport. To demonstrate  
218 this in as clean a fashion as possible, we use targeted experiments with the GMAO FV3  
219 system to show that the degradations in simulated transport when using the remapping  
220 option from the GFDL FV3 core resemble changes that occur when using a lower order

**Table 1. GMAO FV3 Core Finite-Volume Remapping Algorithm:** The remapping algorithm examined in this study controls how individual material surfaces are vertically interpolated back to standard pressure levels. Employing linear – and to a lesser extent – quadratic interpolation in Step 5 produces stratospheric transport characteristics that are more consistent with the most recent GEOS model configurations (green and blue lines, Figure 1), whereas a cubic interpolation is more consistent with older configurations (red and black lines, Figure 1) and with MERRA-2. Here T, U, V, Q,  $C_p$ , K and  $\Phi$  correspond to temperature, zonal wind, meridional wind, specific humidity, specific heat of air at constant pressure and kinetic and potential energy, respectively.

Step	REMAP Procedure (MERRA-2, GMAO FV3 Core)
1	Fit PPM functions to input layer-mean U, V, Q and tracers
2	Calculate PPM to output layer edges
3	Integrate PPM functions between output layer edges to produce new layer-mean U, V, Q and tracers
4	Calculate $TE = C_p T + K + \Phi$ at input mid-layer pressures
5	Calculate TE at output mid-layer pressures using cubic interpolation and a-posteriori integral conservation
6	Construct “remapped” T via $T = (TE - K - \Phi)/C_p$

(quadratic vs. cubic) interpolation scheme in Step 5 of the remapping algorithm (Table 1). This otherwise innocuous change in interpolation order in turn highlights the large sensitivities in transport that can occur even *within* a given finite-volume numerical scheme, in our case resulting in large differences in lower stratospheric upwelling and a  $\sim 30\%$  reduction in the simulated mean age of air.

## 2.2 Model Experiments

### 2.2.1 AMIP vs. EMIP

We begin our analysis by interpreting the results shown in Figure 1, which are all based on historical AMIPs that were performed at the same cubed sphere C180 (approximately half-degree) horizontal resolution. As they represent more official model versions they serve as an important motivation for the experiments that follow. However, there are numerous (potentially compensating) development changes between these model versions which renders it nearly impossible to cleanly identify drivers of differences in their simulated transport.

To this end, in order to investigate the drivers of the differences in Figure 1 we perform targeted modeling experiments aimed at disentangling the influence of individual model development changes on stratospheric transport properties (Table 2). In order to evaluate impacts on transport climate statistics, we consider both a set of climatological AMIP (rows 1-4) as well as so-called “EMIP” (rows 5-7) experiments.

In particular, we carry out 30-year-long AMIP simulations at C180 resolution which we use to infer the climate characteristics of the different model configurations. The “EMIP”

242 experiments – ensembles of 3-month-long integrations initialized on approximately Novem-  
 243 ber 15 of each year between 1985 and 2015 – are also used to infer impacts on simulated  
 244 transport climate. As they are more computationally efficient than AMIPs since all 30  
 245 3-month integrations may be run in parallel, they are performed at both C180 and C360  
 246 (approximately quarter-degree) resolutions in order to examine the sensitivity of our re-  
 247 sults to changes in horizontal resolution.

248 As shown in Appendix B, comparisons of the December-January-February (DJF)  
 249 vertical profile of tropical upwelling show excellent agreement between EMIP and AMIP  
 250 integrations carried out using the same model configuration (Appendix Figure B1). This  
 251 somewhat incidental result represents, to the best of our knowledge, the first time that  
 252 EMIP-based statistics have been shown to converge well to those from AMIPs for the  
 253 stratospheric metrics considered in this study. This suggests that EMIPs, relative to AMIPs,  
 254 may be used to provide a computationally more efficient initial assessment of the impacts  
 255 of model changes on the stratospheric circulation.

### 256 **2.2.2 Model Development Changes**

257 Moving next to the precise model development changes examined, we begin by defin-  
 258 ing a control experiment (CTRL; Table 2, row 1), which best corresponds to the blue  
 259 line shown in Figure 1. Then we define three new AMIP experiments, which aim to dis-  
 260 tinguish between the age-of-air changes resulting from changes in radiation versus changes  
 261 in the handling of the REMAP algorithm that occurred in the transition from **MERRA-**  
 262 **2 to the current FV3 cores** (Section 3.2.1).

263 Specifically, these include experiments in which we a) replace RRTMG with the ra-  
 264 diation from Chou and Suarez (1994) (CSRAD; Table 2, row 2), b) replace the current  
 265 FV3 REMAP approach with the settings used in the GMAO **FV3** core when running  
 266 MERRA-2 (i.e., M2REMAP; Table 2, row 3) and c) combine these two changes (CSRAD+M2REMAP;  
 267 Table 2, row 4).

268 As we show in Section 3.2.1, the M2REMAP experiment produces the largest changes  
 269 in age-of-air, compared to the experiment in which only the radiation is altered. Inter-  
 270 preting this result, however, is not straightforward since there are several differences in  
 271 the implementation of the remapping algorithm between the **GMAO and GFDL FV3**  
 272 cores that are interdependent and, thus, difficult to isolate cleanly. To this end, in or-  
 273 der to simplify the problem we focus the remainder of our investigation (Section 3.2.2)  
 274 on examining a clean set of EMIP experiments that are all performed using the **GMAO**  
 275 **FV3 core remapping settings** and that differ from each other only in terms of the order  
 276 of the interpolation that is used to calculate TE at the mid-layer pressure levels (Step  
 277 5, Table 1). More precisely, we compare configurations using a linear (LINEAR; Table  
 278 2, row 5), quadratic (QUADRATIC; Table 2, row 6) and cubic interpolation (CUBIC;  
 279 Table 2, row 7) scheme, with the latter corresponding to the approach that was used in  
 280 MERRA-2. To assess the robustness of our findings to changes in horizontal resolution,  
 281 all three sensitivity experiments are run at both C180 and C360 resolutions.

282 These three numerical schemes are derived from the generic Lagrangian interpo-  
 283 lation equation:

$$284 \mathcal{Q}(\mathcal{P}) = \sum_{k'} a_{k+k'} \mathcal{Q}_{k+k'} \quad k' = 0, \pm 1, \pm 2, \dots \quad (1)$$

285 where  $\mathcal{P}$  represents the target output location in  $\ln(p)$  and  $\mathcal{Q}_{k+k'}$  denotes the surround-  
 286 ing grid-point values at input locations. The coefficients  $a_{k+k'}$  are derived through Tay-  
 287 lor Series expansions using non-uniform grid spacing given by:

$$288 a_{k+k'} = \frac{\prod_m (\mathcal{P}_{k+m} - \mathcal{P})}{\prod_m (\mathcal{P}_{k+m} - \mathcal{P}_{k+k'})} \quad m = 0, \pm 1, \pm 2, \dots \quad m \neq k' \quad (2)$$



**Table 2. GEOS Model Experiments:** Targeted GEOS model experiments based off a control experiment (row 1) were carried out to identify the influence of radiation (row 2) and changes in the remapping algorithm used since MERRA-2 (row 3), as well as their combined influence (row 4). The influence of the remapping algorithm changes is then interpreted using a simpler set of sensitivity experiments, performed using the **GMAO FV3** core, in which only the order of the interpolation scheme used to calculate TE at output mid-layer pressure levels is altered (rows 5-7). Experiments in rows 1-4 are 30-year-long AMIPs run at C180 resolution, whereas rows 5-7 refer to 30-member 3-month-long (DJF) EMIP experiments. Both AMIPs and EMIPs are used for climate statistic evaluation (see Appendix B for more on the correspondence between the two). EMIP experiments are run at both C180 and C360 horizontal resolutions.

Experiment Name	Configuration	Experiment Type	Hor. Resolution
CTRL	Control, FV3 Core	AMIP (30 yrs.)	C180
CSRAD	Chou-Suarez (1994) Radiation (RAD)	AMIP (30 yrs.)	C180
M2REMAP	GMAO <b>FV3</b> Core (cubic)	AMIP (30 yrs.)	C180
CSRAD+M2REMAP	Chou-Suarez (1994) RAD GMAO <b>FV3</b> Core (cubic)	AMIP (30 yrs.)	C180
LINEAR	GMAO <b>FV3</b> Core (linear)	EMIP (30 mem.)	C180, C360
QUADRATIC	GMAO <b>FV3</b> Core (quadratic)	EMIP (30 mem)	C180, C360
CUBIC	GMAO <b>FV3</b> Core (cubic)	EMIP (30 mem)	C180, C360

287 For the three schemes, the grid points used are: LINEAR ( $k, k-1$ ), QUADRATIC ( $k+$   
288  $1, k, k-1$ ), and CUBIC ( $k+1, k, k-1, k-2$ ). In all cases, the grid points are chosen  
289 such that the target location resides between layers  $k$  and  $k-1$ .

290 Note that, while the LINEAR and QUADRATIC experiments do not actually cor-  
291 respond to any of the model versions shown in Figure 1, they highlight the large sensi-  
292 tivity of the mean age to changes in the interpolation scheme that may otherwise seem  
293 innocuous. They also provide further evidence of the strong influence of changes in trop-  
294 ical lower stratospheric upwelling strength on the stratospheric mean age in GEOS.

295 Finally, in all experiments using the MERRA-2 remapping approach (i.e., M2REMAP,  
296 CSRAD+M2REMAP, LINEAR, QUADRATIC, CUBIC) additional modifications to the  
297 divergence damping coefficients were used so as to best ensure consistency with what was  
298 used in MERRA-2. Specifically, these include changes to the number of layers for ver-  
299 tical subgrid mixing, the coefficient for barotropic mode damping, the use of 2<sup>nd</sup> vs. 6<sup>th</sup>  
300 order divergence damping and the strength of the divergence damping coefficients.

## 301 2.3 Analysis Approach

### 302 2.3.1 Transport Diagnostics

303 To diagnose the transport circulation we focus primarily on the age-of-air (Hall and  
304 Plumb (1994)). This is inferred from an idealized global “clock” or ideal age tracer ( $\Gamma$ )  
305 (Thiele and Sarmiento (1990)) that is defined with respect to the bottom model level  
306 as follows: initially, the ideal age tracer is set to zero throughout the troposphere and  
307 thereafter held to zero over the entire Earth’s surface, subject to a constant aging of 1  
308 year/year throughout the atmosphere. We present here the statistically stationary (equi-  
309 librated) value of  $\Gamma(r)$ , which is equal to the average time since the air at a location  $r$   
310 in the stratosphere last contacted the Earth’s surface. In addition to the mean age, we

also show results from an idealized e90 tracer that is uniformly emitted over the entire surface layer and decays exponentially at a rate of  $90 \text{ days}^{-1}$  such that concentrations greater than 125 ppb and less than 50 ppb tend to reside in the lower troposphere and stratosphere, respectively (Prather et al. (2011)). As this tracer features strong near-tropopause gradients and takes significantly less time to equilibrate, compared to the mean age, it is useful for evaluating stratosphere-troposphere-exchange and transport within the upper troposphere/lower stratosphere (Abalos et al. (2017, 2020); Orbe et al. (2020)).

Both the mean age and e90 tracers were integrated in all of the AMIP experiments shown in Figure 1 and listed in Table 2 (rows 1-4), which were run using the same idealized passive tracer package described in Orbe et al. (2017). Note that the mean age tracer was not integrated in the EMIP experiments given its much longer characteristic timescale in the stratosphere ( $\sim 3\text{-}5$  years). As such, the EMIP simulations, which do not exceed one year, are not appropriate for evaluating the time-integrated transport characteristics reflected in the age-of-air.

In addition to carrying the idealized tracers, two of the experiments shown in Figure 1 were also run with full interactive chemistry and correspond to the two CCMI (Phase 1 and Phase 2) integrations (red and green lines, Figure 1). Both simulations employ the same Global Modeling Initiative (GMI) chemical mechanism (Strahan et al. (2013)) and are therefore useful in evaluating the impact of age differences on real trace gas distributions. In particular, as shown in Section 3.1 results from these experiments show significant imprints of the age-of-air changes on nitrous oxide ( $\text{N}_2\text{O}$ ) and methane ( $\text{CH}_4$ ).

### 2.3.2 Circulation Diagnostics

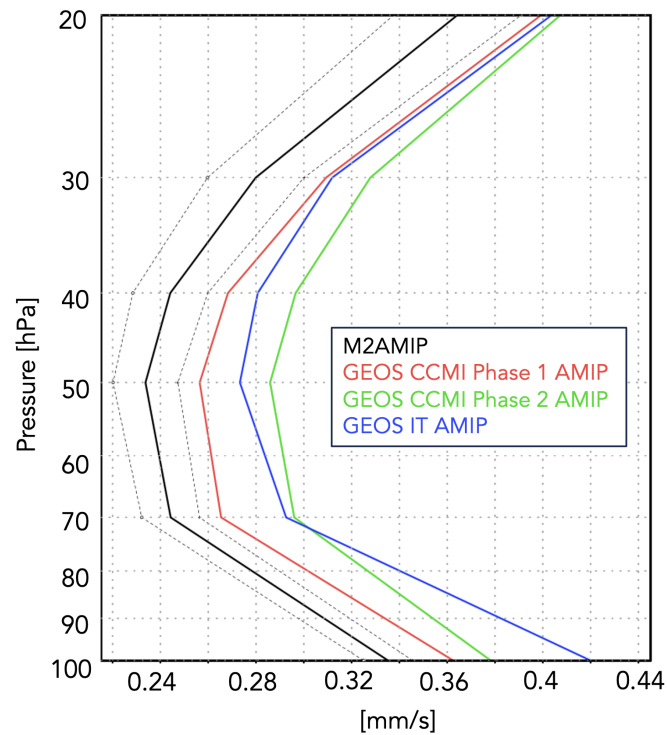
As we show in Section 3, the changes in age-of-air across the different model versions are strongly tethered to changes in the advective component of the circulation, which we quantify using the Transformed Eulerian Mean (TEM) estimate of the Lagrangian transport of mass by the circulation. Thus, in addition to more standard Eulerian metrics of the circulation (e.g., zonal winds and temperatures), we focus on the vertical component of the TEM residual velocity, defined as  $\bar{w}^* = \bar{w} + \frac{\partial(\psi \cos\phi)}{\partial \phi \partial p}$ , where  $\psi = \bar{v}'\theta' / \frac{\partial \theta}{\partial p}$  is the eddy stream function,  $\theta$  refers to potential temperature,  $a$  is the Earth's radius and overbars and primes denote zonal means and deviations therefrom, respectively (Andrews et al. (1987)). In addition, we interpret the behavior in  $w^*$  using the Eliassen-Palm flux divergence ( $\nabla \cdot \mathbf{F}$ ), whose horizontal ( $\mathbf{F}(\phi)$ ) and vertical ( $\mathbf{F}(p)$ ) components are respectively defined as  $\mathbf{F}(\phi) = \cos\phi [\frac{\partial u}{\partial p} \psi - \overline{u'v'}]$  and  $\mathbf{F}(p) = \cos\phi [f - \frac{\partial u \cos\phi}{\partial \phi \partial p}] \psi - \overline{u'\omega'}$ .

## 2.4 Observations and Reanalyses

While our focus is on interpreting and understanding the different model configurations, we incorporate observations to provide context when possible, although we do not present an exhaustive evaluation of the model's transport characteristics (for that see earlier studies including Orbe et al. (2017, 2018)). However, as the tracers are not directly integrated in MERRA-2 (with the exception of ozone), we compare against independent observational estimates. For the mean age we first compare simulated meridional age profiles at 50 hPa with values derived from in situ aircraft measurements of carbon dioxide ( $\text{CO}_2$ ), averaged in 2.5 degree latitude bins over the altitude range 19.5 to 21.5 km (Boering et al. (1996), see also Figure 5 in Hall et al. (1999)).

We also briefly evaluate impacts of transport biases on the simulated trace gas distributions for the CCMI Phase 1 and 2 experiments. The simulated fields of methane ( $\text{CH}_4$ ) are compared with the climatologies derived for 1991–2002 from the Halogen Occultation Experiment (HALOE) on board the Upper Atmosphere Research Satellite (UARS) (Grooß and Russell III (2005)). Comparisons of simulated nitrous oxide ( $\text{N}_2\text{O}$ ) are made against climatologies derived from the Microwave Limb Sounder (MLS) on the Earth Ob-

### DJF Climatological Mean Tropical Upwelling ( $w^*$ )



**Figure 2.** The DJF 1985-1994 climatological mean vertical residual mean velocity,  $w^*$ , averaged at each level between the turnaround latitudes for GEOS free-running AMIP simulations using the model configurations corresponding to the CCM1 Phase 1 (red) and Phase 2 (green) submissions and to GEOS-IT (blue). M2AMIP is shown in black, with black dashed lines denoting  $\pm 1$  standard deviation.

360 serving System (EOS) Aura satellite. Climatologies over the same period (2005–2015)  
 361 are used to evaluate both the model and the observations. We use the 190-GHz retrieval  
 362 from Version 4.2 because the 640-GHz data set ends in summer 2013 due to the failure  
 363 of the N<sub>2</sub>O primary band.

364 For the circulation diagnostics nearly all comparisons are made relative to the MERRA-  
 365 2 data assimilation (DAS) reanalysis product, noting that comparisons against ERA-5  
 366 (not shown) reveal a similar picture. One exception, however, is the vertical component  
 367 of the TEM circulation ( $w^*$ ), which shows some differences in vertical structure between  
 368 the MERRA-2 DAS and a 30-member ensemble of (free-running) AMIP integrations pro-  
 369 duced using the MERRA-2 model, hereafter referred to as M2AMIP (Collow et al., 2017)(Ap-  
 370 pendix Figure A1, right). As the free-running model results shown in Figure 1 show more  
 371 consistency with the vertical profile of M2AMIP, not MERRA-2, we compare  $w^*$  in all  
 372 free-running GEOS experiments with M2AMIP, noting that for non-derived measures  
 373 (i.e., winds, temperatures), the raw MERRA-2 output is used.

374 The differences in  $w^*$  between M2AMIP and the MERRA-2 DAS may reflect the  
 375 influence of temperature increments in the DAS (MERRA-2) which can drive spurious  
 376 vertical transport in assimilated products (Weaver et al., 1993; Orbe et al., 2017). In par-  
 377 ticular, Weaver et al. (1993) showed that the imbalance between the thermal and veloc-  
 378 ity fields at the time an observation is ingested during the assimilation cycle can excite  
 379 unwanted inertial-gravity wave modes that manifest strongly in the residual vertical winds.  
 380 While this impact of the increments may explain the differences in  $w^*$ , we reserve a more-  
 381 in depth examination for future work as our focus is on the transport characteristics of  
 382 the free-running GEOS system.

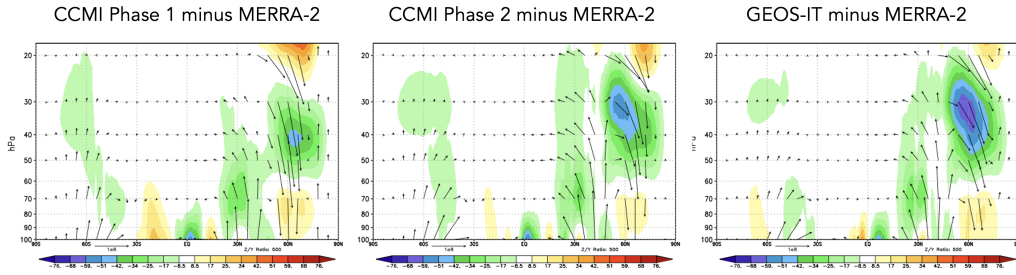
### 383 **3 Results**

#### 384 **3.1 Reduction of Stratospheric Mean Age in GEOS Models** 385 **Since MERRA-2**

386 We begin by interpreting the reduction in mean age exhibited in more recent model  
 387 versions in terms of changes in the strength of upwelling in the tropical lower stratosphere.  
 388 In particular, the reductions in  $\Gamma$  (Figure 1) are consistent with increases in the strength  
 389 of lower stratospheric tropical upwelling, with  $w^*$  becoming progressively stronger in more  
 390 recent model versions, relative to MERRA-2 (Figure 2). Note that these increases in  $w^*$   
 391 across model versions exceed those due to internal variability alone (black dotted lines,  
 392 Fig. 2). Furthermore, while the increases in  $w^*$  occur throughout the stratosphere, we  
 393 focus on the changes occurring between 70 and 100 hPa as these are most relevant to  
 394 determining the tropical upward mass flux and associated strength of the mean overturn-  
 395 ing circulation.

396 Though perhaps naive, the relationship between lower stratospheric upwelling and  
 397 the mean age suggested by comparing Figure 1 and Figure 2 is consistent with the long-  
 398 term behavior of  $\Gamma$  inferred from both historical and projected future climate simulations  
 399 (Butchart et al. (2010); Abalos et al. (2021)). A strong relationship between the strength  
 400 of lower stratospheric ascent and the mean age was also shown to hold in the CCMVal  
 401 models (see Fig. 5.20 in J. Neu et al. (2010)). Nevertheless, it is important to note that  
 402 a clear relationship between  $w^*$  and  $\Gamma$  is not a priori expected, as the age-of-air is also  
 403 known to be very sensitive to mixing, which may be important in interpreting differences  
 404 among the CCM1 Phase 1 models (Dietmüller et al. (2018)).

405 The differences in  $w^*$  highlighted in Figure 2 are associated with enhanced Eliassen-  
 406 Palm flux convergence over NH midlatitudes (Figure 3). Increased wave convergence is  
 407 evident not only within the subtropical lower stratosphere ( $< 30^\circ\text{N}$ , 50-100 hPa) but also  
 408 over higher latitudes and altitudes ( $\sim 40^\circ\text{-}70^\circ\text{N}$ , 20-50 hPa). The fact that differences

DJF Climatological Mean Eliassen-Palm Flux Divergence ( $\nabla \cdot \mathbf{F}$ )

**Figure 3.** Colors show anomalies in the DJF climatological mean Eliassen-Palm (EP) flux divergence between the CCMI Phase 1 (left), CCMI Phase 2 (middle) and GEOS-IT AMIP (right) model versions, relative to MERRA-2. Arrows denote anomalies in the vertical and meridional EP flux vectors (relative to MERRA-2).

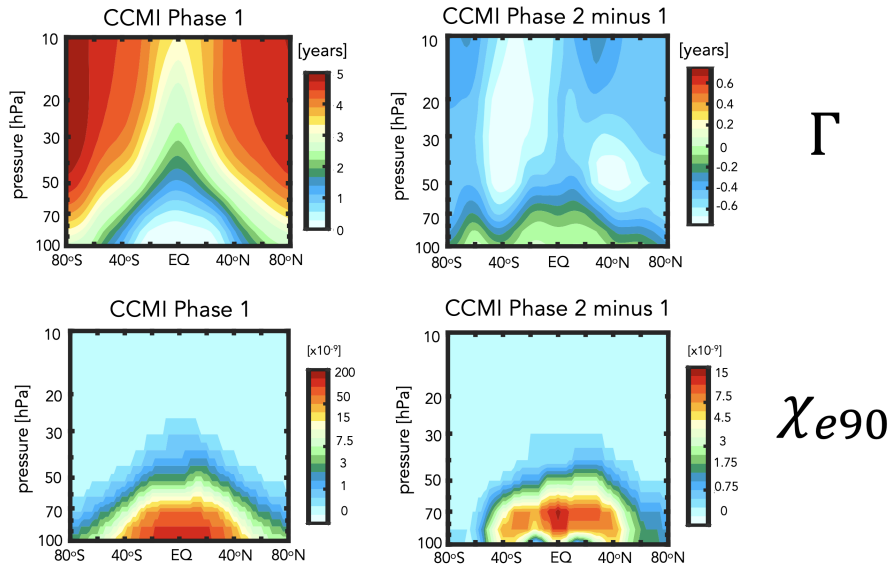
409 in extratropical wave convergence imprint on tropical upwelling is consistent with our  
 410 understanding of the so-called “downward control” principle (Haynes et al. (1991)).

411 In particular, the strength of the residual mean streamfunction ( $\Psi^*$ ) is, via down-  
 412 ward control, directly related to the vertically integrated eddy-induced total zonal force  
 413 above that level and has contributions both from the (resolved wave) Eliassen-Palm flux  
 414 divergence (Figure 3) as well as the gravity wave drag scheme’s parameterized waves (not  
 415 shown). The tropical upward mass flux – defined as  $\Psi_{\max}^* - \Psi_{\min}^*$  evaluated at the turnaround  
 416 latitudes (e.g. Rosenlof (1995)) – is therefore directly dependent on the wave forcing aloft.

417 One subtlety to note is that the wave convergence changes shown in Figure 3 oc-  
 418 cur at high latitudes and are directly associated with downwelling over the polar region.  
 419 It is then via mass balance that anomalously strong downwelling associated with enhanced  
 420 flux convergences must be accompanied by enhanced upwelling in the tropics. This in-  
 421 direct impact of higher latitude wave drag reflects an “extratropical pumping” mecha-  
 422 nism (Holton et al., 1995), which is illustrated more clearly in Section 3.2.2 in the con-  
 423 text of the LINEAR, QUADRATIC and CUBIC experiments.

424 While the reduction in  $\Gamma$  (Figure 1) of  $\sim 30\%$  at 50 hPa is significant, it is neither  
 425 clear if this change is representative of other altitudes within the stratosphere nor how  
 426 this age bias imprints on real chemical species. To this end, we begin by comparing the  
 427 full latitude-pressure distribution of changes in  $\Gamma$  and another passive tracer (e90) (Fig-  
 428 ure 4) between the CCMI Phase 1 and Phase 2 model configurations (red and green lines,  
 429 Figure 1). In particular, we find that the changes in both passive tracers – large reduc-  
 430 tions in  $\Gamma$  within both hemispheres (Fig. 4, top right) and increased values of e90 within  
 431 the lower stratosphere (Fig. 4, bottom right) – are reflective of an overall increase in the  
 432 strength of the transport circulation. This is highlighted in the CCMI Phase 2 – 1 model  
 433 differences for the passive tracer distributions (Fig. 4, right panels) which are shown in  
 434 the absence of robust observational constraints of  $\Gamma$  at higher altitudes (or any obser-  
 435 vational constraints for e90, for that matter). The reduced/increased stratospheric bur-  
 436 dens of the age and e90 tracers are consistent with stronger upwelling in the CCMI Phase  
 437 2 model configuration (Figure 2).

438 While the observational constraints on  $\Gamma$  presented in Figure 1 and the departure  
 439 of  $w^*$  away from MERRA-2 suggest that transport properties of the newer model con-  
 440 figurations are moving in the wrong direction, it is relevant to ask whether or not the



**Figure 4.** The climatological mean (2000-2010) distribution of the mean age-of-air ( $\Gamma$ ) (left, top) and e90 idealized tracers (left, bottom) for the CCMI Phase 1 model configuration. Climatological differences between the CCMI Phase 2 and Phase 1 model configurations are shown in the right panels. Note that a nonlinear colorbar has been used in the e90 subplots.

441 trace gas satellite measurements also support this conclusion. Indeed, comparisons with  
 442 observations show larger biases in  $N_2O$  (Fig. 5, top panels) and  $CH_4$  (Fig. 5, bottom pan-  
 443 els), increasing from 10% to 30% in the CCMI Phase 2 model configuration, depending  
 444 on the species. Recall that the same chemistry mechanism is used in both CCMI Phase  
 445 1 and 2 simulations.

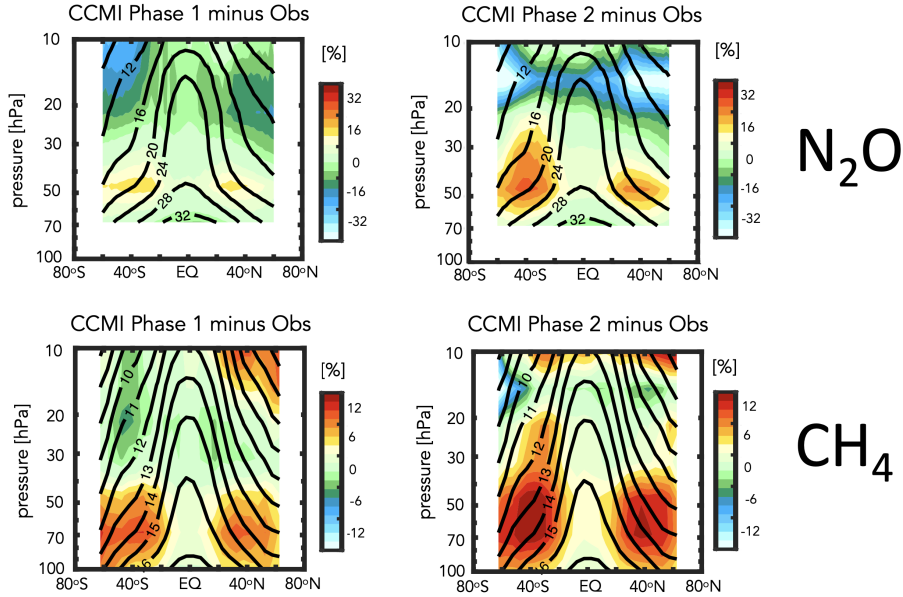
446 The patterns of the trace gases biases are generally consistent with the biases in  
 447 the mean age (Fig. 4). This comports with well-known correlations between the mean  
 448 age and stratospheric trace gases, reinforcing the fact that model transport inaccuracies  
 449 can significantly affect simulations of important long-lived chemical species in the strato-  
 450 sphere (Hall et al. (1999)).

451 The fact that the mean age changes have a significant imprint on the simulated trace  
 452 gases is consequential for the GEOS-R21C system. However, the configurations shown  
 453 in Fig. 1-5 differ in many respects (physics, resolution, radiation, FV remapping algo-  
 454 rithm) and it is difficult to meaningfully interpret what is driving the changes in  $w^*$  (and  
 455 the tracers). We therefore move next to the targeted model experiments (Table 2) in or-  
 456 der to interpret the model development steps that resulted in these transport circula-  
 457 tion changes.

## 458 **3.2 Identifying Drivers of Upwelling and Tracer Changes Since** 459 **MERRA-2**

### 460 **3.2.1 Radiation versus REMAP Algorithm**

461 As discussed in Section 2, among the model changes that were made since MERRA-  
 462 2, the changes in radiation and the FV remapping algorithm are most likely to directly  
 463 have impacted the stratospheric circulation. We therefore begin by assessing which of  
 464 these changes dominates the decreases in  $\Gamma$  shown in Figure 1.



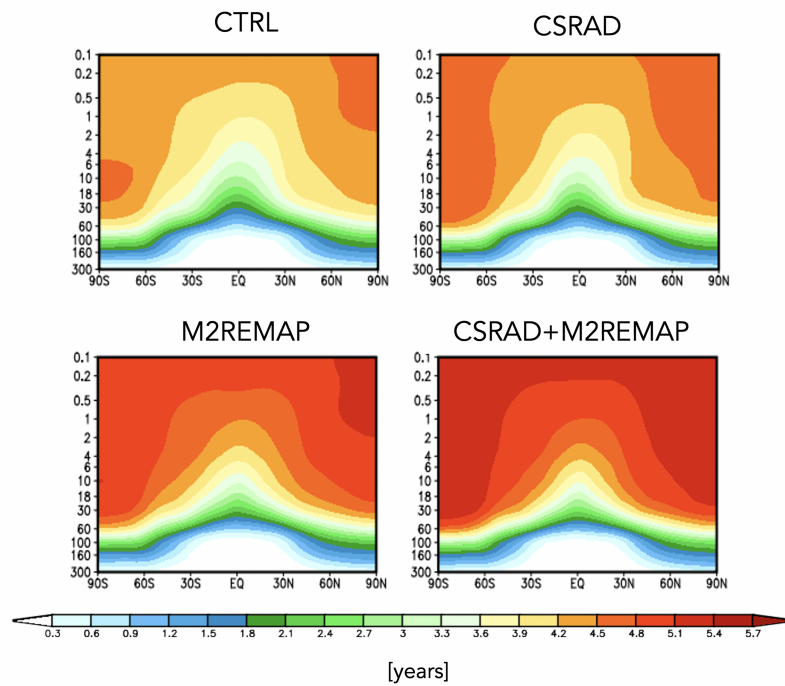
**Figure 5.** Colors shown anomalies in the simulated distributions of nitrous oxide ( $\text{N}_2\text{O}$ ) (top) and methane ( $\text{CH}_4$ ) (bottom), relative to the MLS and HALOE observed values, respectively, for the CCMI Phase 1 (left) and Phase 2 (right) GEOS model configurations. Climatological mean observed values are shown in the black contours.

465 Figure 6 shows the distribution of  $\Gamma$  for experiments in which the longwave, short-  
 466 wave, and REMAP updates since MERRA-2 have successively been undone. Relative  
 467 to the control experiment (CTRL; Table 2, row 1), the transition back to Chou and Suarez  
 468 (1994) in the shortwave and Chou (1990, 1992) in the longwave results in an increase in  
 469 the mean age of  $\sim 0.5$  years throughout the stratosphere (CSRAD; Table 2, row 2). Though  
 470 significant, this change in  $\Gamma$  is smaller than the change that results from applying the  
 471 remapping approach used in MERRA-2 (M2REMAP; Table 2; row 3), in which the mean  
 472 age increases by  $\sim 1$  year. The combined impacts of both changes (CSRAD+M2REMAP;  
 473 Table 2 row 4) is roughly linear, with age values of  $\sim 5.5$  years over high latitudes at 50  
 474 hPa, consistent with the values simulated by the GEOS-CTM MERRA-2 integration (black  
 475 line, Figure 1) and with the CCMI Phase-1 version of the model (red line, Figure 1).

476 Next we ask if the behavior of  $\Gamma$  exhibited in Figure 6 can be interpreted in terms  
 477 of changes in the strength of lower stratospheric tropical upwelling and extratropical wave  
 478 convergence, as our previous analysis of the CCMI experiments suggested. Indeed, Fig-  
 479 ure 7 shows that values of upwelling decrease in the CSRAD and M2REMAP experiments,  
 480 relative to the CTRL integration. The increase in upwelling resulting from both changes  
 481 (CSRAD+M2REMAP) is still larger, consistent with the larger age decreases in that ex-  
 482 periment. This change in the behavior of  $w^*$  within the tropical stratosphere can be in-  
 483 terpreted in terms of changes in the Eliassen Palm flux convergence over NH midlati-  
 484 tudes (not shown), which features smaller values in the CSRAD, M2REMAP (and CSRAD+  
 485 MSREMAP) experiments. Note that our examination of the changes in  $w^*$  are derived  
 486 from EMIP integrations, which we showed previously converge (for DJF) to the statis-  
 487 tics derived from corresponding AMIP experiments.

488 **It is important to note that, while the reduced values of  $w^*$  in the M2REMAP ex-**  
 489 **periment (Fig. 7), represent an improvement, relative to the CTRL integration, these**

### Annual Climatological Mean Stratospheric Mean Age ( $\Gamma$ )



**Figure 6.** Colors show the simulated 2000-2010 climatological annual mean distributions of the mean age-of-air ( $\Gamma$ ) for the CTRL (top left; Table 2, row 1), CSRAD (top right; Table 2, row 2), M2REMAP (bottom left; Table 2, row 3) and combined CSRAD+M2REMAP (bottom right; Table 2, row 4) experiments.



changes are, in isolation, accompanied by a cold bias ( $\sim 4\text{-}6\text{K}$ ) in the tropical lower stratosphere (Fig. 8a). Unlike the upwelling changes, this represents a degradation in model skill, increasing the temperature bias in that region, relative to MERRA-2 (Fig. 8 b,c). While this temperature bias is concerning in isolation, it is decoupled from the changes in residual mean upwelling which, rather, are more sensitive to the changes in *extratropical* wave convergences (described further in the next section). Furthermore, the increased temperature bias in the M2REMAP experiment appears to be tethered to other updates that were made in the CTRL model (specifically the radiation progression from Chou-Suarez to RRTMG). In particular, the M2AMIP ensemble, which also employs the GMAO cubic remapping option, features a much smaller temperature bias in the tropical lower stratosphere, relative to MERRA-2 ( $\sim 0\text{-}2\text{K}$ ) (Fig. 8d). This suggests that the amplified temperature bias moving from the CTRL to M2REMAP experiments needs to be interpreted in the context of other model development changes that were introduced, particularly in the radiation scheme and its coupling to convection.

### 3.2.2 FV REMAP Algorithm: Sensitivity of Climate Statistics

Having shown in the previous section that the largest changes in the mean age were realized through the differences in implementation of the remapping algorithm between the GMAO FV3 core used in MERRA-2 and in current FV3 core configurations, we now investigate further the sensitivity of the transport circulation to the choice of remapping interpolation scheme. In particular, we compare simulations run with the GMAO FV3 remapping settings in which total energy is calculated at new mid-layer pressures using cubic, quadratic and linear interpolation prior to the aposterior integral conservation (Table 2, rows 5-7). In addition, in this section we seek to understand how the changes in the Eliassen-Palm flux convergence over NH midlatitudes arise via analysis of the large-scale wind structure.

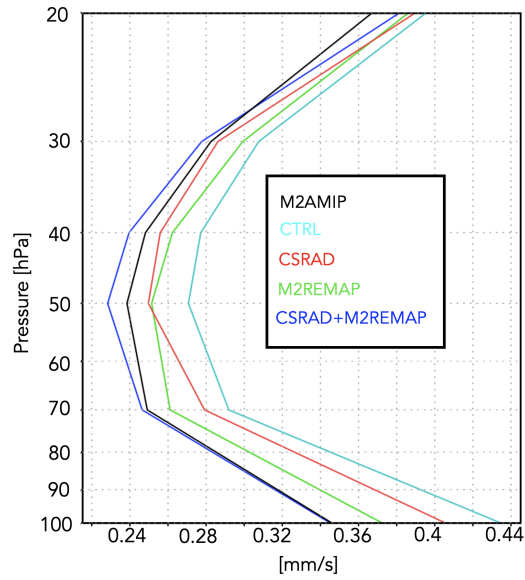
Figure 9 (left panel) shows a clear sensitivity in tropical upwelling to the choice of interpolation scheme, with  $w^*$  progressively increasing in strength moving from the CUBIC to QUADRATIC to LINEAR schemes. This sensitivity is robust across horizontal resolutions as the same suite of experiments performed at C360 exhibit the same sensitivity (Fig. 9, right panel). While no current model version actually employs a linear scheme, this suite of experiments highlights the strong sensitivity to choice of interpolation scheme within the remapping algorithm; to the best of our knowledge, this result has not been reported in the literature.

Furthermore, as we show next, this clean set of experiments allow us to inquire mechanistically into the processes that are driving the changes in wave convergence over midlatitudes, unencumbered by differences in horizontal resolution, physics, etc.

Interestingly, the increases in  $w^*$  moving from the CUBIC to QUADRATIC schemes not only manifests in free-running AMIP simulations, but also in AMIP simulations in which GEOS is constrained (or replayed, following Orbe et al. (2017)) to MERRA-2 meteorological fields, using the MERRA-2 GMAO cubic interpolation (Fig. 10, blue line) and the GFDL FV3 core remapping approach (Fig. 10, red line). While there is a general increase in  $w^*$  in the former, however, both simulations, lie within the range of MERRA2-DAS, suggesting that replay does act to ameliorate some of the upwelling biases manifest in the underlying unconstrained models. While not the focus of this study, this impact of the remapping approach on simulations run in replay mode will be examined further in future work.

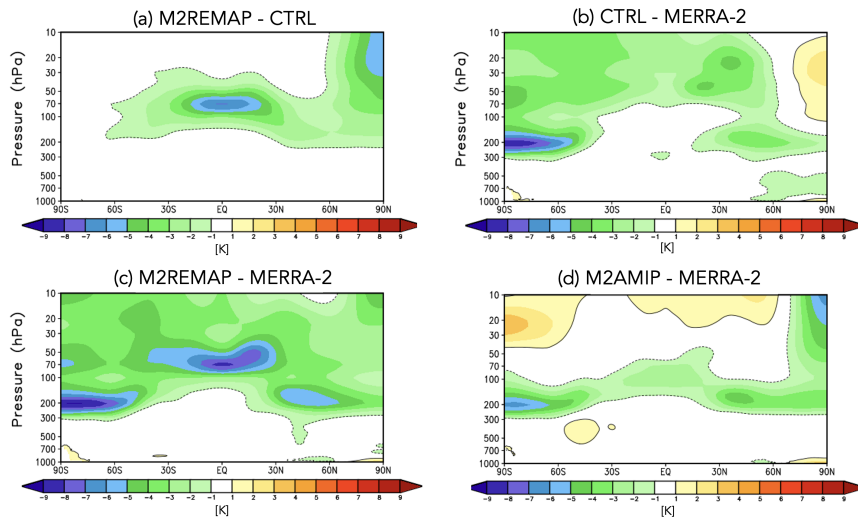
Consistent with our expectations based on the analysis of the previous experiments, the drivers of the changes in  $w^*$  are related to increased wave convergence moving from the CUBIC to QUADRATIC to LINEAR schemes (Figure 11). Over extratropical latitudes, the zonal force associated with this enhanced wave convergence is associated with enhanced downwelling at high latitudes that, through mass balance, is accompanied by

### DJF Climatological Mean Tropical Upwelling ( $w^*$ )



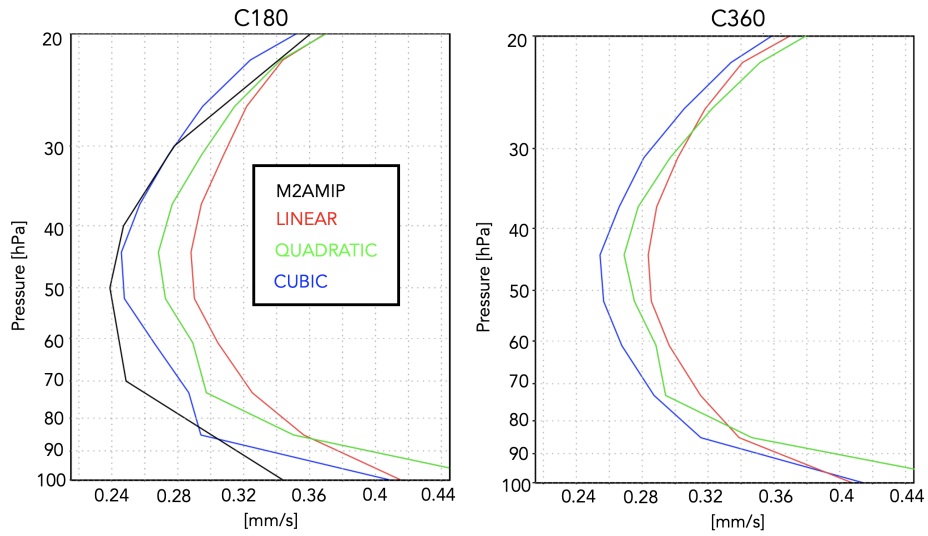
**Figure 7.** The DJF 1985-2015 climatological mean vertical residual mean velocity,  $w^*$ , averaged at each level between the turnaround latitudes for the CTRL (cyan line; Table 2, row 1), CSRAD (red line; Table 2, row 2), M2REMAP (green line; Table 2, row 3) and combined CSRAD+M2REMAP (blue line; Table 2, row 4) experiments. M2AMIP is shown in black.

### DJF Climatological Zonal Mean Temperature Anomalies



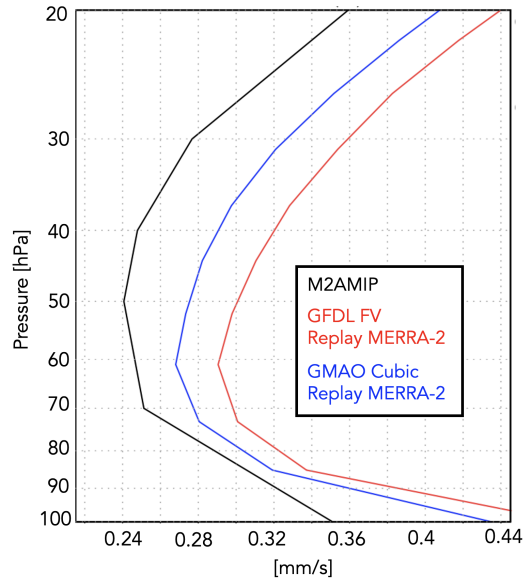
**Figure 8.** Anomalies in DJF 1985-2015 climatological mean zonal mean temperatures: M2REMAP - CTRL (a), CTRL - MERRA-2 (b), M2REMAP - MERRA-2 (c) and M2AMIP - MERRA-2 (d).

### DJF Climatological Mean Upwelling ( $w^*$ )

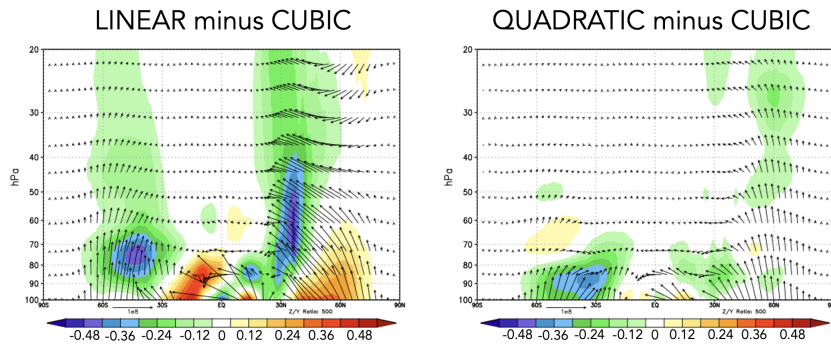


**Figure 9.** The DJF 1985-2015 climatological mean vertical residual mean velocity,  $w^*$ , averaged at each level between the turnaround latitudes for the LINEAR (red line; Table 2, row 5), QUADRATIC (green line; Table 2, row 6) and CUBIC (blue line; Table 2, row 7) experiments. M2AMIP is shown in black. Results from C180 and C360 EMIP experiments are shown in the left and right panels, respectively.

### DJF Climatological Mean Tropical Upwelling ( $w^*$ )



**Figure 10.** The DJF 2005-2015 climatological mean vertical residual mean velocity,  $w^*$ , averaged at each level between the turnaround latitudes for two GEOS replay AMIP simulations constrained with MERRA-2 meteorological fields using remapping approaches from the MERRA-2 GMAO **FV3** cubic core (blue line) and the GFDL **FV3** core (red line). MERRA-2 DAS is shown in black.

DJF Climatological Mean Eliassen-Palm Flux Divergence ( $\nabla \cdot \mathbf{F}$ )

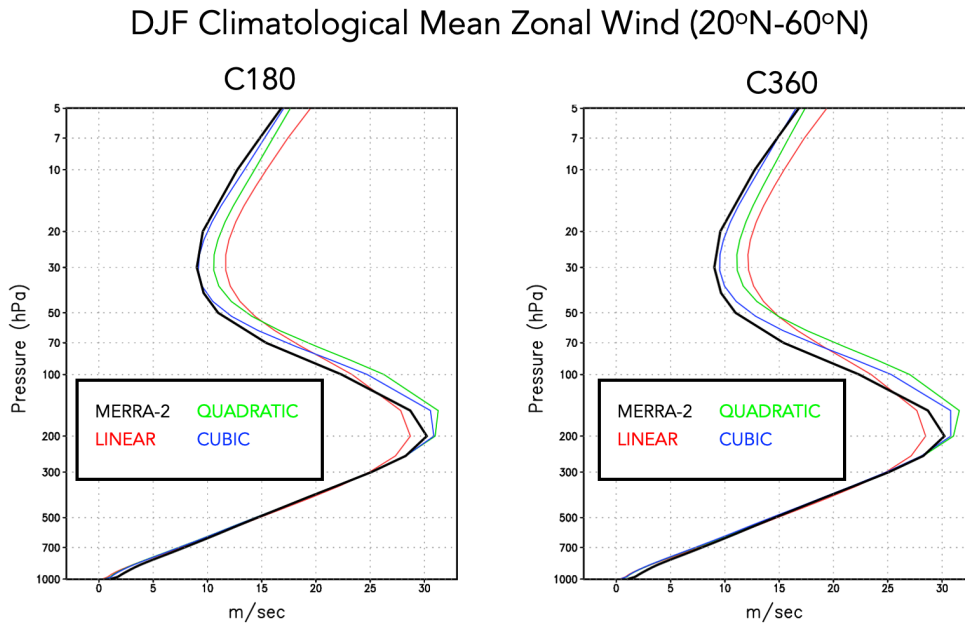
**Figure 11.** Colors shown anomalies in the DJF 1985-2015 climatological mean Eliassen-Palm (EP) flux divergence in the LINEAR (left) and QUADRATIC (right) experiments, relative to the CUBIC model experiment. Arrows denote anomalies in the vertical and meridional EP flux vectors.

541 enhanced upwelling in the tropics. This indirect impact of higher latitude wave drag is  
 542 evident in Appendix Figure C1, which shows stronger upwelling/downwelling in the LIN-  
 543 EAR and QUADRATIC experiments over the tropics/polar region.

544 Next we exploit the fact that these experiments only differ with respect to the in-  
 545 terpolation scheme in order to inquire further into the drivers of the wave convergence  
 546 changes. To this end, **Figure 12** compares profiles of the zonal mean zonal wind between  
 547 the CUBIC, QUADRATIC and LINEAR experiments, averaged over the region of en-  
 548 hanced wave convergence (i.e. 20°N-60°N). The experiments featuring stronger wave con-  
 549 vergence (LINEAR and QUADRATIC) are also simulations with stronger zonal winds,  
 550 relative to MERRA-2, especially above 70 hPa. This change in winds occurs at both C180  
 551 (**Fig. 12**, left panel) and C360 (**Fig. 12**, right panel) resolutions.

552 Structurally, the increase in zonal wind strength over northern extratropical mid-  
 553 latitudes is reflective of a poleward shift in the zonal winds as the critical latitude, i.e.  
 554 where the zonal wind is zero, shifts northward in the QUADRATIC and, especially, LIN-  
 555 EAR integrations, relative to the CUBIC experiment (**Figure 13**). Since stationary waves  
 556 only propagate in westerly zonal flow, the latitude where zonal flow is zero acts a bound-  
 557 ary for wave propagation (Hardiman et al. (2014)). As a result, this shift in critical lat-  
 558 itude results in enhanced wave propagation and convergence over middle and high lat-  
 559 itudes.

560 **Figures 12 and 13** highlight how the changes in zonal winds in the LINEAR and  
 561 QUADRATIC experiments reflect a degradation in model skill, relative to MERRA-2,  
 562 throughout the entire stratosphere. The changes in upwelling, mean age, chemical trace  
 563 gases and zonal winds thus provide a coherent and self-consistent picture suggestive of  
 564 a degradation in the representation of the stratospheric circulation since MERRA-2. That  
 565 is, an increased bias in the stratospheric northern zonal winds are, via their influence on  
 566 wave convergence, compromising changes in the strength of the mean meridional over-  
 567 turning circulation and its impact on composition. It is interesting to note that the wind  
 568 biases also extend into the troposphere and show degraded skill relative to MERRA-2  
 569 in the LINEAR and QUADRATIC experiments (**Figure 13**). Examination of other fields



**Figure 12.** Vertical profiles of the DJF 1985-2015 climatological mean zonal mean zonal winds in the LINEAR (red), QUADRATIC (green) and CUBIC (blue) experiments, averaged between 20°N and 60°N. MERRA-2 is shown in the black line. Results for both C180 (left) and C360 (right) experiments are shown.

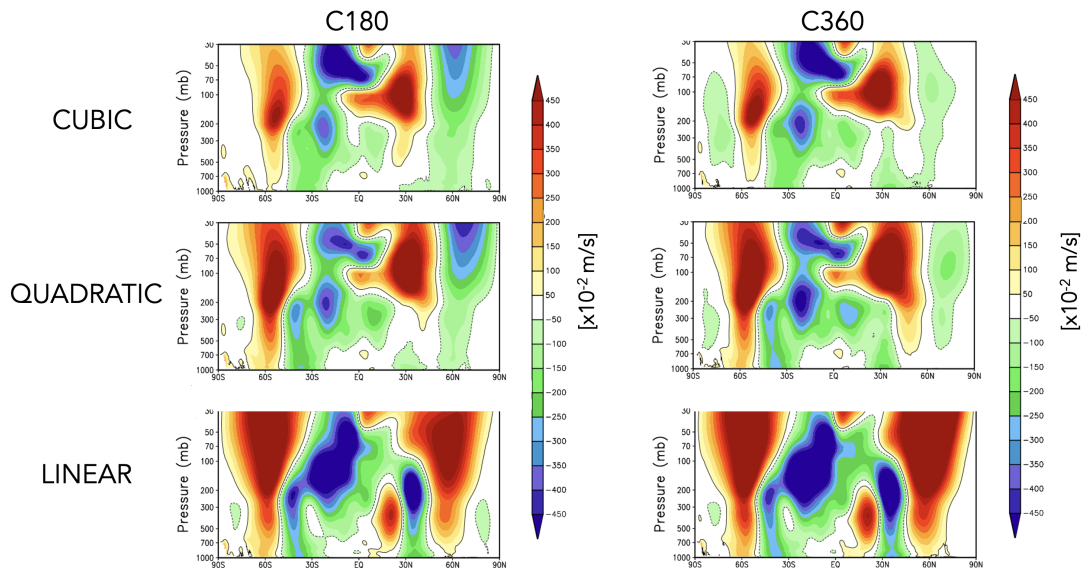
570 (i.e. tropopause biases, Appendix Figure D1) present somewhat more of a nuanced story  
 571 that depends on other changes that made to the model since MERRA-2, as shown in Fig-  
 572 ure 8 and discussed in Section 3.2.1. The improvements in the zonal winds, however, are  
 573 most relevant for setting the upwelling characteristics within the tropical lower strato-  
 574 sphere via their influence on wave propagation into that region.

575 Finally, to better understand why these impacts on the winds have such a conse-  
 576 quence for the wave convergence properties within the stratosphere, next we examine the  
 577 zonal structure of these biases in the middle stratosphere (Figure 14). This reveals that  
 578 the enhanced winds in the LINEAR (and, to a lesser extent, QUADRATIC) integrations  
 579 are concentrated over the North Pacific at both C180 (Fig. 14, top) and C360 (Fig. 14,  
 580 bottom) resolutions (a similar picture emerges within the troposphere, not shown). As  
 581 this region is the primary region dominating the stationary component of the upward  
 582 flux of vertical wave activity (Plumb (1985), see their Figure 4) it is perhaps not surpris-  
 583 ing that this region is having a profound impact on the mean overturning circulation.  
 584 Again, as with the zonal mean wind changes, the increases in wind strength over the North  
 585 Pacific represent degraded model skill relative to MERRA-2. Note that comparisons with  
 586 ERA-5 reveal a similar bias (not shown).

## 587 4 Conclusions

588 Here we have performed an analysis aimed at understanding differences in the repre-  
 589 sentation of the stratospheric circulation in recent candidate systems for GEOS-R21C,  
 590 relative to older versions of GEOS similar to the model used to produce MERRA-2. Us-  
 591 ing targeted experiments oriented at disentangling various model development updates,  
 592 we have identified a key role played by changes in the implementation of the remapping

## DJF Climatological Zonal Mean Zonal Wind Anomalies Relative to MERRA-2



**Figure 13.** Colors shown anomalies in the DJF 1985-2015 climatological mean zonal mean zonal winds in the CUBIC (top), QUADRATIC (middle) and LINEAR (bottom) experiments, relative to MERRA-2. Results for both C180 (left) and C360 (right) experiments are shown.

593 algorithm within the model’s finite-volume dynamical core. Our key results are as fol-  
 594 lows:

595 #1. The stratospheric mean age-of-air in GEOS is sensitive to the degree of the  
 596 interpolation scheme that is used to calculate layer-mean values of total energy,  $U$ ,  $V$  and  
 597 tracers. Different treatment of the vertical remapping algorithm result in mid-stratospheric  
 598 (50 hPa) age-of-air differences of  $\sim 1$  year over high latitudes, or about 30% climatolog-  
 599 ical mean values.

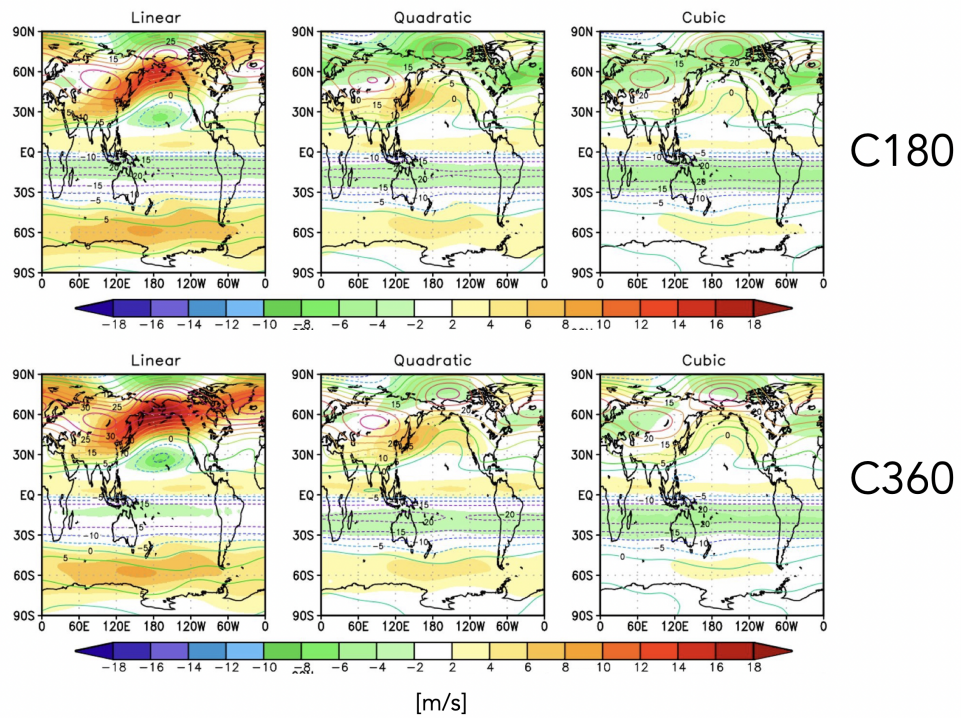
600 #2. The increased age-of-air biases in more recent GEOS configurations are reflected  
 601 in the increased biases in simulated trace gases, including  $\text{CH}_4$  and  $\text{N}_2\text{O}$ .

602 #3. The age-of-air sensitivities reflect, to first order, changes in the strength of trop-  
 603 ical upwelling associated with the Brewer-Dobson circulation which are in turn are driven  
 604 by changes in EP flux convergence over northern midlatitudes. Changes in wave conver-  
 605 gence reflect shifts in (critical lines of) wave propagation that originate in the troposphere  
 606 over the Pacific Ocean, a region of strong upward wave activity.

607 #4. The degradation of upwelling statistics manifest in AMIPs, also translate to  
 608 degradations in configurations of GEOS in which the meteorological fields are constrained  
 609 or “replayed” to MERRA-2.

610 An important caveat should be mentioned in relation to Conclusion 3 listed above,  
 611 which highlights the leading role played by *extratropical* wave convergences in the be-  
 612 havior of tropical upwelling in the model. While this is consistent with our understand-  
 613 ing that the Brewer-Dobson Circulation is, to first order, determined by extratropical  
 614 wave flux convergences, in practice this means that other undesirable biases may emerge  
 615 in the tropics (i.e., tropical temperatures, Figure 8). Our analysis of the M2AMIP en-  
 616 semble, however, suggests that this tropical temperature bias is not a necessary conse-

### DJF Climatological 30 hPa Zonal Wind Anomalies Relative to MERRA-2



**Figure 14.** Colors shown anomalies in the DJF 1985-2015 climatological mean zonal winds at 30 hPa in the CUBIC (right), QUADRATIC (middle) and LINEAR (left) experiments, relative to MERRA-2. Results for both C180 (top) and C360 (bottom) experiments are shown.

617 quence of the GMAO FV3 remapping scheme but, rather, emerges in tandem with other  
 618 changes that were made in the model. Nonetheless, Figure 8 suggests that improvements  
 619 in transport associated with the GMAO remapping scheme must be carefully weighed  
 620 in the broader context of other model development choices.

621 Interestingly, preliminary analysis suggests that our findings may also translate to  
 622 replay configurations of GEOS (Figure 10), although the effect is muted, relative to free-  
 623 running configurations. As a rigorous evaluation of the stratospheric circulation in re-  
 624 play and DAS configurations is beyond the scope of the current study, future work will  
 625 therefore focus on assessing the extent to which the free-running model biases reported  
 626 here are expressed when the model is run in data assimilation mode. It also bears em-  
 627 phasizing that our findings do have immediate implications for the (free-running) sub-  
 628 seasonal forecast and coupled chemistry climate applications of the GEOS model cur-  
 629 rently in operation.

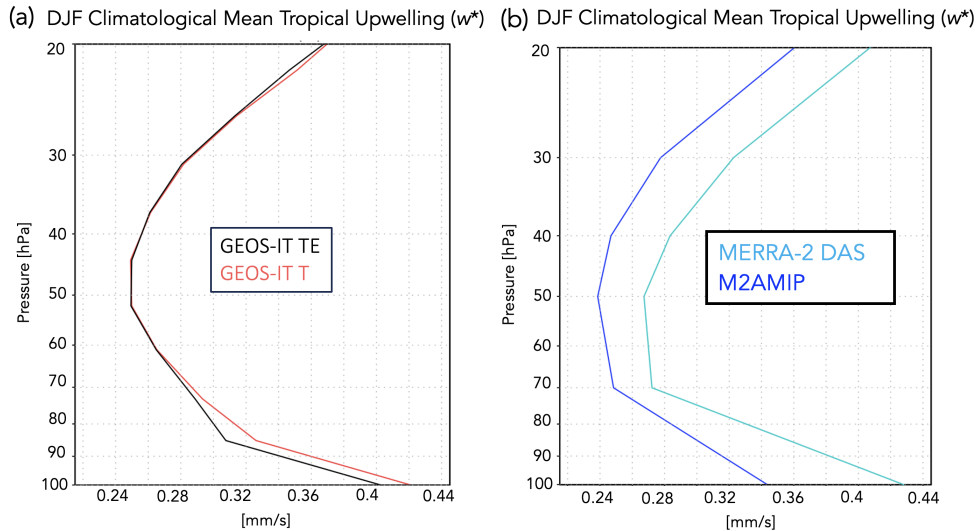
630 In addition to its implications for GEOS, our results more generally highlight the  
 631 key role played by model numerics in transport (e.g., Rood (1987)). The sensitivities in  
 632 the age-of-air documented herein are also consistent in spirit with the findings in Gupta  
 633 et al. (2020) who showed significant age differences occurring between spectral versus  
 634 finite-volume numerics. Our results, however, suggest that there remain large sensitiv-  
 635 ities even within a given (FV) dynamical core.

636 Looking forward, our findings support and build on the recommendation proposed  
 637 in Gupta et al. (2020) for the construction of dynamical core benchmark tests aimed at  
 638 determining how underlying AGCM numerics impact climatological transport proper-  
 639 ties. In particular, in addition to the age-of-air, the authors propose a range of strato-  
 640 spheric circulation diagnostics that should be evaluated including the zonal mean zonal  
 641 winds, eddy temperature variance and zonal spectra of eddy kinetic energy. Our anal-  
 642 ysis reveals an important role to be played by the climatological zonal mean wind struc-  
 643 ture as it impacts wave convergence over midlatitudes; we therefore also recommend ex-  
 644 plicit consideration of the Eliassen Palm flux convergence and tropical upwelling ( $w^*$ )  
 645 fields as they may be crucial for interpreting age-of-air changes.

646 One somewhat incidental – but practical - result from our analysis is that the statis-  
 647 tics of  $\nabla \cdot \mathbf{F}$  and  $w^*$  are well approximated by ensembles of so-called EMIP integrations.  
 648 As these are substantially easier to run than AMIPs these could provide a “first pass”  
 649 when evaluating new proposed model development changes, without the immediate need  
 650 to integrate AMIP-style experiments. We emphasize, however, that this statement should  
 651 only apply to a first stage in model development as the age-of-air will reflect the time  
 652 integrated impacts of both advection and mixing.

653 Finally, we conclude by noting that, while we have focused on sensitivities within  
 654 the FV remapping algorithm, our results have highlighted important sensitivities to changes  
 655 in radiation and, to a lesser extent, changes in parameterized convection. Though not  
 656 the dominant drivers of the age-of-air changes identified here, the former could poten-  
 657 tially influence the age both directly through changes in thermal structure and indirectly  
 658 by modifying wave propagation and/or generation in the troposphere.





**Figure A1.** The DJF 1985-2015 climatological mean vertical residual mean velocity,  $w^*$ , averaged at each level between the turnaround latitudes compared between two experiments remapping to temperature (T) (red) versus total energy (TE) (black) (a) and between MERRA-2 DAS (cyan) and the M2AMIP ensemble (blue) (b). The underlying model code is consistent with the Version 5.29.4 GEOS-IT model.

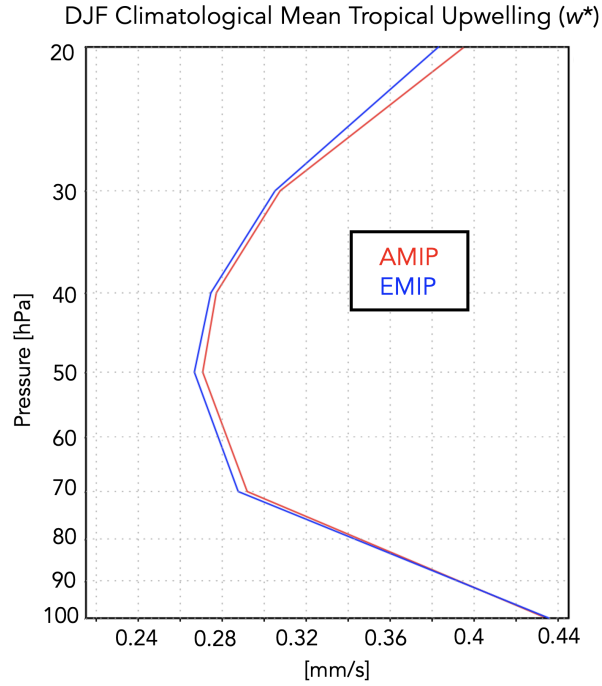
## 659 Appendix A Sensitivities in Calculation of TEM Upwelling

660 There are various aspects of the calculation of the TEM circulation that warrant  
 661 further comment. First, whereas the modeling experiments listed in Table 2 (rows 5-7)  
 662 focus on the sensitivity of Step 5 within the **GMAO FV3** core remapping algorithm to  
 663 the choice of interpolation scheme, another difference between the **GMAO** and **GFDL**  
 664 **FV3** core remapping approaches is the use of TE versus T, respectively. To test the im-  
 665 pact of this difference, we ran a new experiment which is identical to the CUBIC exper-  
 666 iment (Table 2, row 7), except that T is remapped from input layer mean pressure lo-  
 667 cations to standard output layer mean locations directly using cubic interpolation (i.e.,  
 668 no computation of TE or a-posteriori energy conservation applied). Appendix Figure A1a  
 669 shows that this has little impact on the strength of tropical upwelling, suggesting that  
 670 the  $w^*$  differences associated with changes in the remapping algorithm are dominated  
 671 by sensitivities to the choice of interpolation scheme, not the use of TE versus T.

672 Second, the vertical component of the TEM circulation ( $w^*$ ) shows some differences  
 673 in vertical structure between MERRA-2 and the 30-member M2AMIP ensemble (Ap-  
 674 pendix Figure A1b). This difference in vertical structure appears to reflect a difference  
 675 between DAS and free-running configurations of the model, since other DAS configura-  
 676 tions share a similar vertical structure (not shown). Given this difference, we ensure as  
 677 apples-to-apples a comparison of simulated TEM velocities by comparing all AMIP re-  
 678 sults to other AMIP experiments.

## 679 Appendix B Correspondence between EMIP and AMIP Upwelling

680 Appendix Figure B1 shows the close correspondence in DJF climatological mean  
 681  $w^*$ , averaged at each level between the turnaround latitudes, from AMIP and EMIP ex-

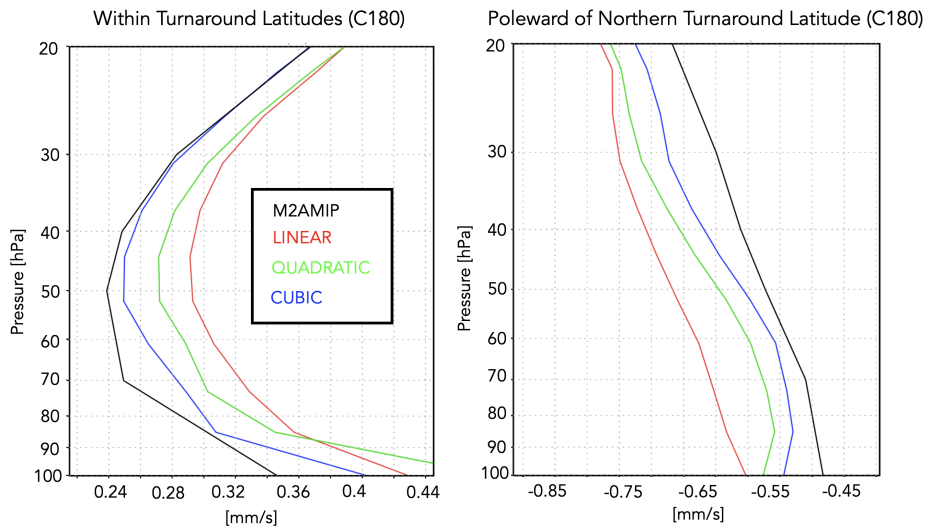


**Figure B1.** The DJF 1985-2015 climatological mean vertical residual mean velocity,  $w^*$ , averaged at each level between the turnaround latitudes for the CTRL experiment (Table 2, row 1). Results based on a 30-year-long AMIP experiment (red line) and a 30-member ensemble of three-month-long EMIP experiments (blue line) are shown.

682 periments using the same model configuration. This good agreement in upwelling is used  
 683 to justify the analysis of the EMIP experiments listed in Table 2 (rows 5-7).

## 684 **Appendix C Changes in Tropical and High Latitude Upwelling**

685 Appendix Figure C1 compares the behavior in residual mean upwelling among the  
 686 LINEAR, QUADRATIC and CUBIC experiments over the latitudes between the (tropi-  
 687 cal) turnaround latitudes (left) and poleward of the northern turnaround latitude (right).  
 688 The ordering among experiments in both regions reflects how increases in downwelling  
 689 at high latitudes are, through mass balance, accompanied by enhanced upwelling in the  
 690 tropics.

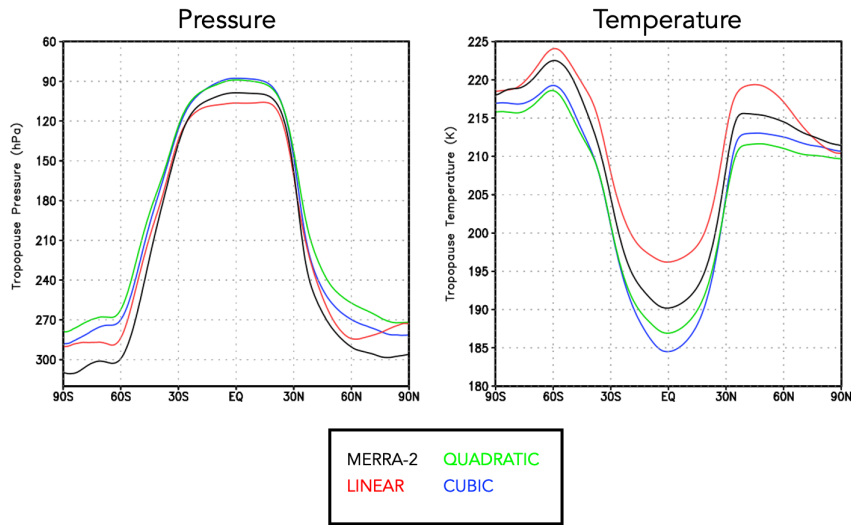
DJF Climatological Mean Upwelling ( $w^*$ )

**Figure C1.** Left: The DJF 1985-2015 climatological mean vertical residual mean velocity,  $w^*$ , averaged at each level between the turnaround latitudes for the LINEAR (red line; Table 2, row 5), QUADRATIC (green line; Table 2, row 6) and CUBIC (blue line; Table 2, row 7) experiments. M2AMIP is shown in black. Right: As in left panel, except averaged over latitudes poleward of the northern turnaround latitude. Results in both panels are shown for C180 experiments.

691 **Appendix D Tropopause Pressure**

692 Appendix Figure D1 compares boreal winter tropopause pressure and temperature  
 693 among the LINEAR, QUADRATIC and CUBIC experiments, relative to MERRA-2.

## DJF Climatological Mean Tropopause



**Figure D1.** The DJF 1985-2015 climatological mean tropopause pressure (left) and temperature (right) in the CUBIC (blue), QUADRATIC (green) and LINEAR (red) experiments. MERRA-2 is shown in black. Results are shown for the C180 experiments.

### 694 Acronyms

- 695 **AMIP** Atmospheric Model Intercomparison Project  
 696 **CH<sub>4</sub>** methane  
 697 **CCMs** chemistry climate models  
 698 **CCMI** Chemistry Climate Modeling Initiative  
 699 **CCMVal** Chemistry Climate Model Validation  
 700 **CO<sub>2</sub>** carbon dioxide  
 701 **CTRL** control  
 702 **CTM** chemistry transport model  
 703 **DAS** Data assimilation  
 704 **DJF** December-January-February  
 705 **EMIP** ensemble AMIP  
 706 **EOS** Earth Observing System  
 707 **EP** Eliassen-Palm  
 708 **FV** finite-volume  
 709 **FP** Forward Processing  
 710 **GEOS** Global Earth Observing System  
 711 **GEOS-R21C** GEOS Retrospective analysis for the 21<sup>st</sup> Century  
 712 **GMI** Global Modeling Initiative  
 713 **HALOE** Halogen Occultation Experiment  
 714 **MERRA-2** Modern-Era Retrospective Analysis for Research and Applications v2  
 715 **MLS** Microwave Limb Sounder  
 716 **N<sub>2</sub>O** nitrous oxide  
 717 **NH** northern hemisphere  
 718 **PPM** piecewise parabolic  
 719 **RRTMG** Rapid Radiative Transfer Model for GCMS

720 **SW** shortwave  
 721 **TE** total energy  
 722 **TEM** Transformed Eulerian Mean  
 723 **UARS** Upper Atmosphere Research Satellite

## 724 **Open Research Section**

725 TBD

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