

1 **Non-monotonic feedback dependence on CO<sub>2</sub> due to a**  
2 **North Atlantic pattern effect**

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9 **Key Points:**

- 10 • Effective climate sensitivity (EffCS) and radiative feedbacks change non-monotonically  
11 with increasing CO<sub>2</sub> concentrations
- 12 • The non-monotonicity is associated with the formation of a sea-surface temperature  
13 cooling pattern over the North Atlantic
- 14 • Results imply an overlooked radiative damping effect on global EffCS from the North  
15 Atlantic warming hole

## Abstract

Effective climate sensitivity (EffCS), commonly estimated from model simulations with abrupt CO<sub>2</sub> quadrupling, has been shown to depend on the level of CO<sub>2</sub> forcing. To understand this dependency systematically, we performed a series of simulations with a range of abrupt CO<sub>2</sub> forcing in two climate models (CESM1-LE and GISS-E2.1-G). Our results indicate that EffCS is a non-monotonic function of the CO<sub>2</sub> forcing, decreasing between 3× and 4×CO<sub>2</sub> in CESM1 (2× and 3×CO<sub>2</sub> in GISS) and increasing at higher CO<sub>2</sub> levels. The minimum EffCS value, caused by anomalously negative radiative feedbacks, arises from sea-surface temperature (SST) cooling in the North Atlantic. This localized North Atlantic cooling pattern is associated with the formation of the North Atlantic Warming Hole, accompanied by the collapse of the Atlantic Meridional Overturning Circulation under CO<sub>2</sub> forcing. Our findings emphasize the importance of understanding changes in North Atlantic SST patterns for constraining near-future and equilibrium global warming.

## Plain Language Summary

Estimates of effective climate sensitivity (EffCS) are complicated by 1) the nonlinear dependence of feedback on temperature and 2) the sea-surface temperature (SST) pattern effect. We find that EffCS and radiative feedbacks change non-monotonically with CO<sub>2</sub> concentrations due to a cooling SST pattern over the North Atlantic associated with the formation of a “North Atlantic Warming Hole” (NAWH). While most previous studies focused on the impact of tropical Pacific SST patterns on EffCS, we here highlight an overlooked damping effect on EffCS from North Atlantic SST cooling across CO<sub>2</sub> levels. Our results imply that understanding and constraining the NAWH under CO<sub>2</sub> forcings is crucial for transient warming projections and EffCS constraints.

## 1 Introduction

Equilibrium climate sensitivity (ECS), the equilibrium global-mean surface air temperature response to a doubling of atmospheric CO<sub>2</sub> relative to pre-industrial (PI) levels, is one of the most important metrics in climate science. The Charney 1979 report estimated a “likely” ECS range of 1.5-4.5K; most recently, a tighter range of ECS values between 2.6-3.9K was established using a Bayesian framework that combines multiple lines of evidence (Sherwood et al., 2020).

When evaluated from climate models, ECS is often approximated with an effective climate sensitivity (EffCS), estimated from 150-year abrupt CO<sub>2</sub> quadrupling simulations within coupled global climate models (GCMs), with an underlying assumption that EffCS remains constant with different CO<sub>2</sub> doublings. However, previous modeling (Bloch-Johnson et al., 2021; Meraner et al., 2013; Mauritsen et al., 2019; Sherwood et al., 2020; Mitevski et al., 2021; Zhu & Poulsen, 2020) and paleoclimate studies (Anagnostou et al., 2016, 2020; Farnsworth et al., 2019; Friedrich et al., 2016; Shaffer et al., 2016; Zhu et al., 2019) have shown that EffCS may not be linear with each successive CO<sub>2</sub> doubling. It tends to increase at higher CO<sub>2</sub> values primarily due to a nonlinear temperature dependence of the radiative feedbacks ( $\lambda$ ), referred to as the state-dependence of feedbacks (Andrews et al., 2015; Bloch-Johnson et al., 2021; Sherwood et al., 2015), with minor contributions from nonlinear CO<sub>2</sub> dependence of radiative forcing (Mitevski et al., 2022).

However, previous attempts to study the state dependence have been limited to CO<sub>2</sub>-doubling scenarios (2×, 4×, 8×CO<sub>2</sub>) (Good et al., 2016; M. Rugenstein et al., 2019), whereas the shared socioeconomic pathway for the highest emission scenarios (SSP5-8.5) projects a transient increase of greenhouse gas forcing up to 8×CO<sub>2</sub> at the year 2300 (Meinshausen et al., 2020) passing through all the intermediate states of  $n\times\text{CO}_2$  ( $n = 2, 3, 4, 5, 6, 7, 8$ ). Moreover, previous studies have been focused on how EffCS and feedbacks vary in response to changes in global-mean temperatures under different CO<sub>2</sub> forcing (Meraner et al., 2013;

65 Caballero & Huber, 2013; Bloch-Johnson et al., 2021), with little attention on how the  
 66 spatial patterns of feedback and local surface warming respond to various CO<sub>2</sub> forcings.  
 67 Hence in this study, we examine the dependence of EffCS on CO<sub>2</sub> levels systematically and  
 68 its connection to the spatial patterns of the climate feedbacks by performing and analyzing  
 69 a hierarchy of GCM experiments with a range of abrupt CO<sub>2</sub> forcings including 2×, 3×,  
 70 4×, 5×, 6×, 7×, and 8×CO<sub>2</sub> relative to PI level (hereafter denoted as abrupt  $n\times\text{CO}_2$   
 71 experiments).

## 72 2 Materials and Methods

### 73 2.1 Models and Experiments

74 We use the original large ensemble version of the Community Earth System Model  
 75 (CESM1-LE). CESM1-LE comprises the Community Atmosphere Model version 5 (CAM5,  
 76 30 vertical levels) and parallel ocean program version 2 (POP2, 60 vertical levels) with  
 77 approximately 1° horizontal resolution in all model components (Kay et al., 2015). Some of  
 78 the results are shown with the GISS-E2.1-G model (Kelley et al., 2020) in Supplementary  
 79 Information. All experiments in this work are with abrupt CO<sub>2</sub> forcing.

80 We perform abrupt  $n\times\text{CO}_2$  experiments with the coupled version of the CESM1-LE  
 81 and GISS-E2.1-G models (coupled runs) for 150 years with 2×, 3×, 4×, 5×, 6×, 7×, and  
 82 8×CO<sub>2</sub> forcing, with all other trace gases, aerosols, ozone concentrations, and solar forcing  
 83 fixed at PI values. The response is defined as the difference between the  $n\times\text{CO}_2$  runs and  
 84 the PI control run. The same experiments were analyzed in (Mitevski et al., 2021, 2022).

85 To estimate the effective radiative forcing (ERF) as per Forster et al. (2016), we per-  
 86 form prescribed pre-industrial SST and sea-ice runs for 30 years for each 2×, 3×, 4×, 5×,  
 87 6×, 7×, and 8×CO<sub>2</sub>. The ERF is then calculated as the global mean net top of the atmo-  
 88 sphere (TOA) net radiation between PI and  $n\times\text{CO}_2$ , and it includes the stratospheric and  
 89 tropospheric adjustments (Sherwood et al., 2015).

90 We also utilize atmosphere-only runs (AGCM) with prescribed monthly SST values  
 91 taken from the 150-year abrupt  $n\times\text{CO}_2$  runs. The prescribed SST values are monthly data  
 92 for 150 years. The CO<sub>2</sub> concentration, ozone concentrations, aerosols, solar forcing, and all  
 93 other trace gases are fixed at pre-industrial values.

In addition to only prescribing SST values from the  $n\times\text{CO}_2$  runs, we also change the  
 SST patterns. We use the pattern from 3×CO<sub>2</sub> in CESM1-LE and then scale the pattern  
 by the global-mean warming amplitude from 4×CO<sub>2</sub> and 5×CO<sub>2</sub>. We do this by

$$\Delta\text{SST}(x, y, t) = \text{SST}_{3\times\text{CO}_2}(x, y, t) - \text{SST}_{\text{PI}}(x, y),$$

$t$  is monthly data from 150 years,  $x$  is longitude, and  $y$  is latitude. Next, we find the pattern  
 $S_p$  as

$$S_p(x, y, t) = \frac{\Delta\text{SST}(x, y, t)}{\overline{\Delta\text{SST}}(t)}$$

where  $\overline{\Delta\text{SST}}$  is the global mean monthly data for 150 years. Then we have

$$\Delta\text{SST}'_{n\times\text{CO}_2}(x, y, t) = S_p \cdot \overline{\text{SST}_{n\times\text{CO}_2}}.$$

And finally

$$\text{SST}_{n\times\text{CO}_2}(x, y, t) = \text{SST}_{\text{PI}}(x, y) + \Delta\text{SST}'_{n\times\text{CO}_2}(x, y, t).$$

94 One caveat here is that we are only changing the SSTs, and holding sea-ice fixed at 3×CO<sub>2</sub>.  
 95 Although sea-ice changes also cause albedo feedback changes, their contribution is much  
 96 smaller than the SST-mediated feedback changes.

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## 2.2 Analysis

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We calculate the effective climate sensitivity EffCS as the x-intercept of regressing the change in net TOA radiation against surface air temperature over the 150 years of the simulations (Gregory et al., 2004; Zelinka et al., 2020). We normalize the EffCS by  $\log_2 n$  for the  $n \times \text{CO}_2$  runs, assuming logarithmic  $\text{CO}_2$  forcing, consistent with Bloch-Johnson et al. (2021).

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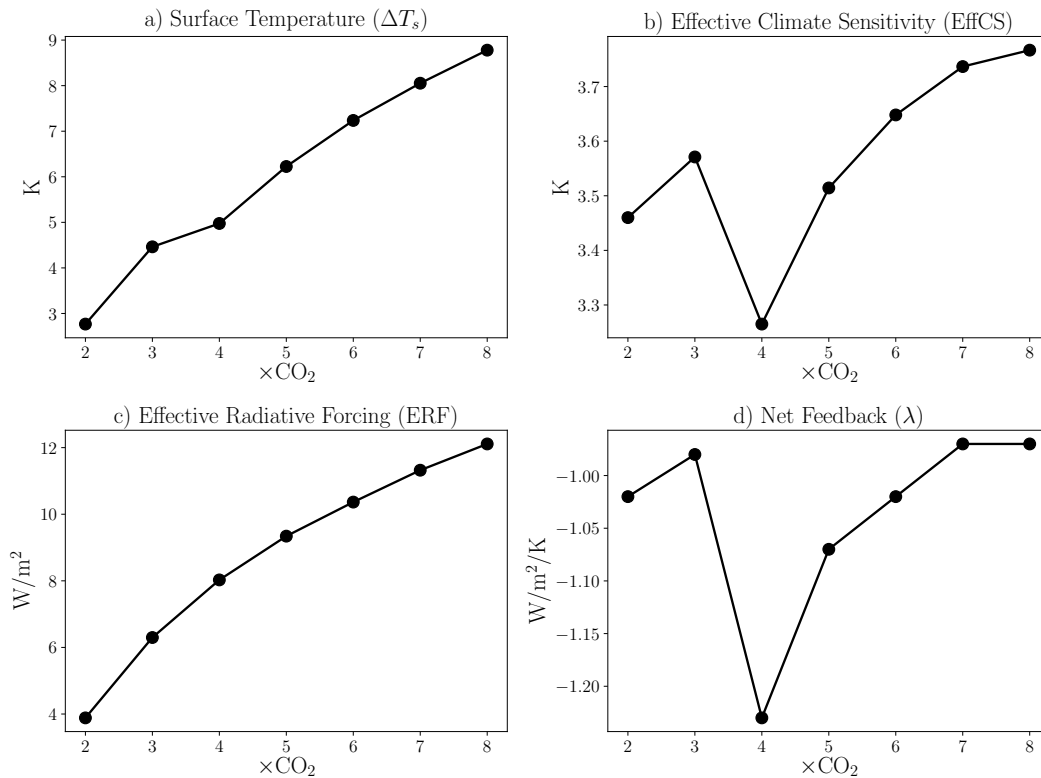
We calculate individual feedbacks with radiative kernels from Pendergrass et al. (2018). For each year, we multiply the spatially resolved kernels by the climate field anomalies of atmospheric temperature  $T$ , water vapor  $q$ , and surface albedo  $\alpha$ . We regress these quantities on the surface temperature response, and the slope of this regression is the feedback. The cloud feedbacks are computed via the residual method (Soden & Held, 2006).

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## 3 Results

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### 3.1 Non-monotonic effective climate sensitivity and radiative feedbacks



**Figure 1.** a) Global mean surface air temperature response ( $\Delta T_s$ ), b) effective climate sensitivity (EffCS), c) effective radiative forcing (ERF) from 30-year fixed sea-surface temperature runs, and d) net feedback parameter ( $\lambda$ ) from the 150-year Gregory regression of abrupt  $n \times \text{CO}_2$  runs.

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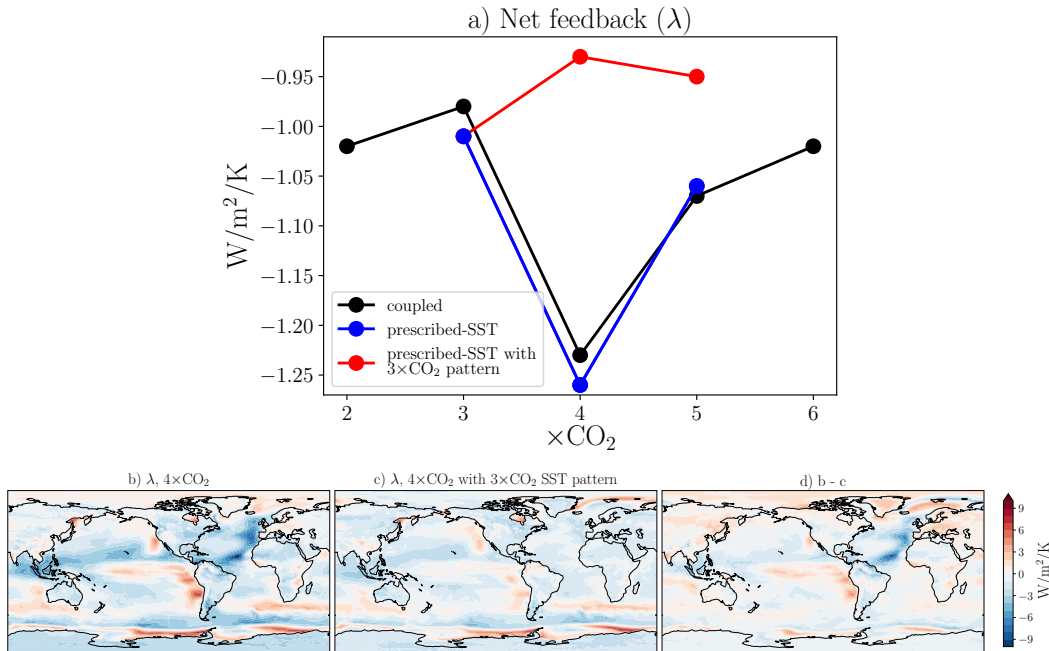
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Results from CESM1-LE show that although the global-mean surface air temperature increases monotonically as  $\text{CO}_2$  increases (Fig. 1a), EffCS changes non-monotonically with  $\text{CO}_2$  levels (Fig. 1b). That is, EffCS decreases between  $3 \times$  and  $4 \times \text{CO}_2$  and then increases between  $4 \times$  and  $5 \times \text{CO}_2$ , and at higher  $\text{CO}_2$  forcing, with a minimum value at  $4 \times \text{CO}_2$ . We find the same non-monotonicity in the GISS-E2.1-G experiments except with a minimum EffCS at  $3 \times \text{CO}_2$  (Fig. S1). In the rest of the paper, we will focus on the CESM1-LE

116 simulations and note that the results hold for the GISS-E2.1-G simulations unless otherwise  
 117 noted.

118 Changes in EffCS, in principle, are governed by changes in effective radiative forcing  
 119 (ERF) and radiative feedbacks ( $\lambda$ ). While the ERF, calculated from an additional 30-year  
 120 fixed sea-surface temperature (SST) runs as per Forster et al. (2016), increases slightly more  
 121 than the logarithm of the  $\text{CO}_2$  concentration at higher  $\text{CO}_2$  levels than  $4\times\text{CO}_2$  (see Mitevski  
 122 et al. (2022) for more detail), it is strongly monotonic with  $\text{CO}_2$  and does not exhibit a minimum  
 123 value (Fig. 1c). On the other hand, the net radiative feedback parameter  $\lambda$  (Fig. 1d),  
 124 calculated from 150-year regressions of top-of-atmosphere (TOA) radiative response against  
 125 surface air temperature change (Zelinka et al., 2020), exhibits a clearly non-monotonic be-  
 126 havior with respect to  $\text{CO}_2$  levels, as for EffCS:  $\lambda$  becomes more negative (more stabilizing)  
 127 between  $3\times$  and  $4\times\text{CO}_2$  and less negative between  $4\times$  and  $5\times\text{CO}_2$ , corresponding to the  
 128 lowest EffCS at  $4\times\text{CO}_2$ . Similar results are also found in the GISS-E2.1-G model experi-  
 129 ments (Fig. S1). These results suggest that EffCS depends not only *nonlinearly* on  $\text{CO}_2$ , as  
 130 found in previous studies (Meraner et al., 2013; Caballero & Huber, 2013; Bloch-Johnson et  
 131 al., 2021), but also *non-monotonically*, and that the non-monotonicity is caused by the ra-  
 132 diative feedbacks in our simulations. Hence the question is: what causes the non-monotonic  
 133 changes in feedbacks?



**Figure 2.** a) Global net feedback parameter  $\lambda$  from coupled runs, AGCM prescribed-SST runs with SSTs from coupled runs, and prescribed-SST runs with  $3\times\text{CO}_2$  pattern, where the  $3\times\text{CO}_2$  SST patterns are scaled with the actual global-mean SST values of  $4\times\text{CO}_2$  and  $5\times\text{CO}_2$ , respectively. Spatial patterns of the local contribution to the global  $\lambda$  at  $4\times\text{CO}_2$  from b) prescribed-SST, c) prescribed-SST with  $3\times\text{CO}_2$  pattern and d) the difference.

### 134 3.2 Non-monotonic $\lambda$ traced to changes in surface warming patterns

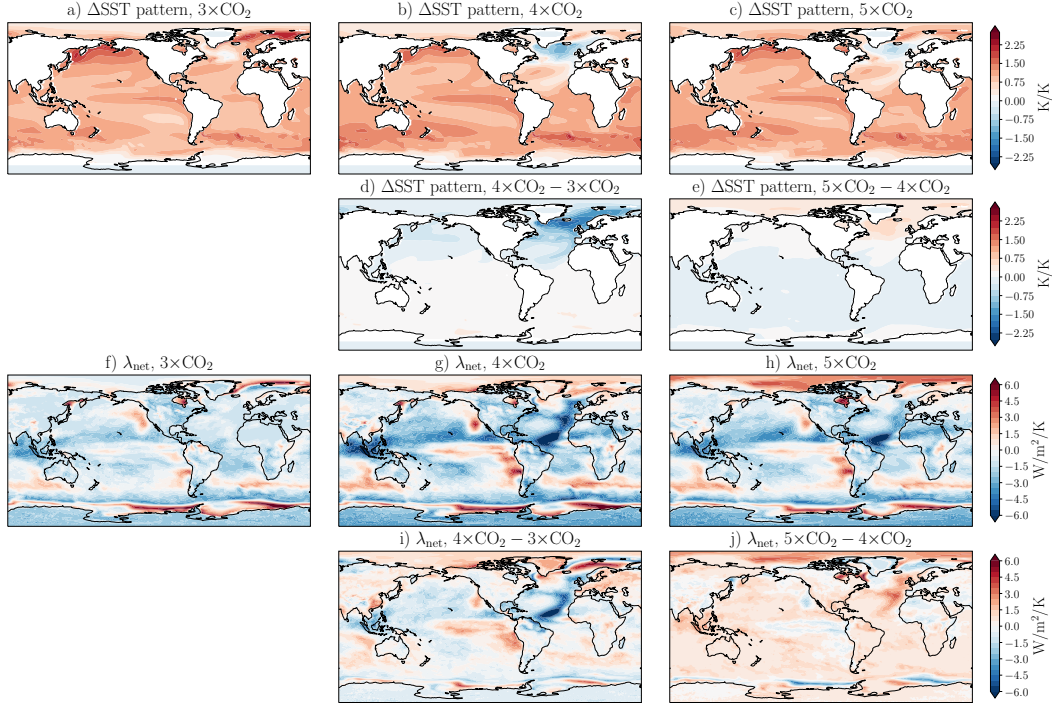
135 We hypothesize two reasons for the non-monotonic changes in  $\lambda$  with  $\text{CO}_2$ :

- 136 1. The non-monotonic dependence in  $\lambda$  may arise from a nonlinear state-dependence of  
 137 the feedbacks. As noted above, previous studies have found that radiative feedbacks  
 138 change nonlinearly with global-mean surface temperature changes (i.e., feedback tem-  
 139 perature dependence), mostly owing to the cloud and water vapor feedbacks (Meraner  
 140 et al., 2013; Caballero & Huber, 2013; Bloch-Johnson et al., 2021). Can the changes  
 141 in global-mean surface temperature across the CO<sub>2</sub> levels in our simulations (Fig. 1a)  
 142 explain the non-monotonic behavior of  $\lambda$  and, therefore, EffCS?
- 143 2. The non-monotonic dependence of  $\lambda$  may arise from a strong dependence of  $\lambda$  on the  
 144 spatial pattern of SSTs. Recent studies have found a close coupling between SST  
 145 patterns and radiative feedbacks in observations and model simulations, the so-called  
 146 “pattern effect” (Zhou et al., 2016; Dong et al., 2019; Sherwood et al., 2020). If the  
 147 SST pattern effect caused the non-monotonic response in  $\lambda$ , then what SST regions  
 148 govern the global and local changes in our feedbacks?

149 To test the hypotheses, we run the atmospheric component of the coupled model  
 150 CESM1-LE (CAM5) with specified SST boundary conditions, in order to examine the im-  
 151 pacts of different surface warming on  $\lambda$ . First, we perform a set of 150-year long CAM5  
 152 simulations where we fix all radiative forcing agents at pre-industrial levels, and prescribe  
 153 the time-varying SSTs produced by the corresponding coupled model  $n \times \text{CO}_2$  simulations.  
 154 In these runs (denoted as “prescribed-SST”), TOA radiative fluxes and surface air tempera-  
 155 ture freely adjust to the underlying SSTs. Although not directly forced by CO<sub>2</sub>, we find that  
 156 the prescribed-SST simulations accurately reproduce the values of  $\lambda$  from the corresponding  
 157 coupled simulations (c.f. blue and black dots in Fig. 2a). This finding, consistent with other  
 158 studies (Haugstad et al., 2017; Zhou et al., 2023), suggests that the dependence of  $\lambda$  on CO<sub>2</sub>  
 159 forcing is primarily shaped by the SSTs induced by the CO<sub>2</sub> forcing and therefore confirms  
 160 the validity of using prescribed-SST simulations to study radiative feedbacks to understand  
 161 the coupled  $n \times \text{CO}_2$  results.

162 Next, we perform another set of prescribed-SST simulations with adjusted SST bound-  
 163 ary conditions. To test hypothesis # 1, i.e., whether  $\lambda$  responds non-monotonically to  
 164 changes in global-mean surface temperatures, we conduct simulations where we scale the  
 165 SST pattern from  $3 \times \text{CO}_2$  by the actual global-mean SST changes in coupled  $4 \times \text{CO}_2$  and  
 166  $5 \times \text{CO}_2$ , respectively. Such that these two runs have the same normalized global SST pat-  
 167 tern (at every monthly time step) as the  $3 \times \text{CO}_2$  run but different global-mean SST values  
 168 (denoted “prescribed-SST with  $3 \times \text{CO}_2$  pattern”). In these experiments, we find that the  $\lambda$   
 169 values do not reproduce those in the coupled & prescribed-SST simulations even though the  
 170 same global-mean SST warming is prescribed (c.f. red and blue dots in Fig. 2a), suggesting  
 171 that the non-monotonic response in  $\lambda$  arises from changes in the spatial pattern of SSTs  
 172 (hypothesis # 2) and not the changes in the global-mean values of SSTs (hypothesis # 1).

173 The above CAM5 prescribed-SST simulations highlight the role of SST patterns in  
 174 driving the non-monotonic response in  $\lambda$ . To understand what regions contribute to this  
 175 non-monotonicity, we show the spatial pattern of  $\lambda$  calculated as the local net TOA radiation  
 176 regressed to global-mean surface air temperature response, shown in Fig. 2b-d. The spatial  
 177 pattern of  $\lambda$  in the  $4 \times \text{CO}_2$  prescribed-SST run ( $4 \times \text{CO}_2$  SST pattern), is shown in Fig. 2b,  
 178 corresponding to the globally averaged  $\lambda$  at  $4 \times \text{CO}_2$  shown by the blue dot in Fig. 2a. The  
 179 spatial pattern of  $\lambda$  in a  $4 \times \text{CO}_2$  run with  $3 \times \text{CO}_2$  SST pattern (red dot in Fig. 2a) is shown  
 180 in Fig. 2c. Taking the difference between Fig. 2b and c (panel d) shows substantially more  
 181 negative feedback in the North Atlantic with the  $4 \times \text{CO}_2$  pattern, and not much change  
 182 when we use the  $3 \times \text{CO}_2$  pattern, indicating that the anomalously low EffCS at  $4 \times \text{CO}_2$  in  
 183 our coupled simulations is primarily associated with an anomalously negative  $\lambda$  in the North  
 184 Atlantic.



**Figure 3.** Maps of SST patterns (calculated as the regression of local temperature changes to global temperature changes for 150 years) in the coupled runs for a)  $3\times\text{CO}_2$ , b)  $4\times\text{CO}_2$ , and c)  $5\times\text{CO}_2$ . The differences between  $4\times$  and  $3\times\text{CO}_2$ , and  $5\times$  and  $4\times\text{CO}_2$ , are shown in d) and e), respectively. Figures f-j) show  $\lambda$  maps for the same  $\text{CO}_2$  experiments.

### 3.3 A local pattern effect from the North Atlantic

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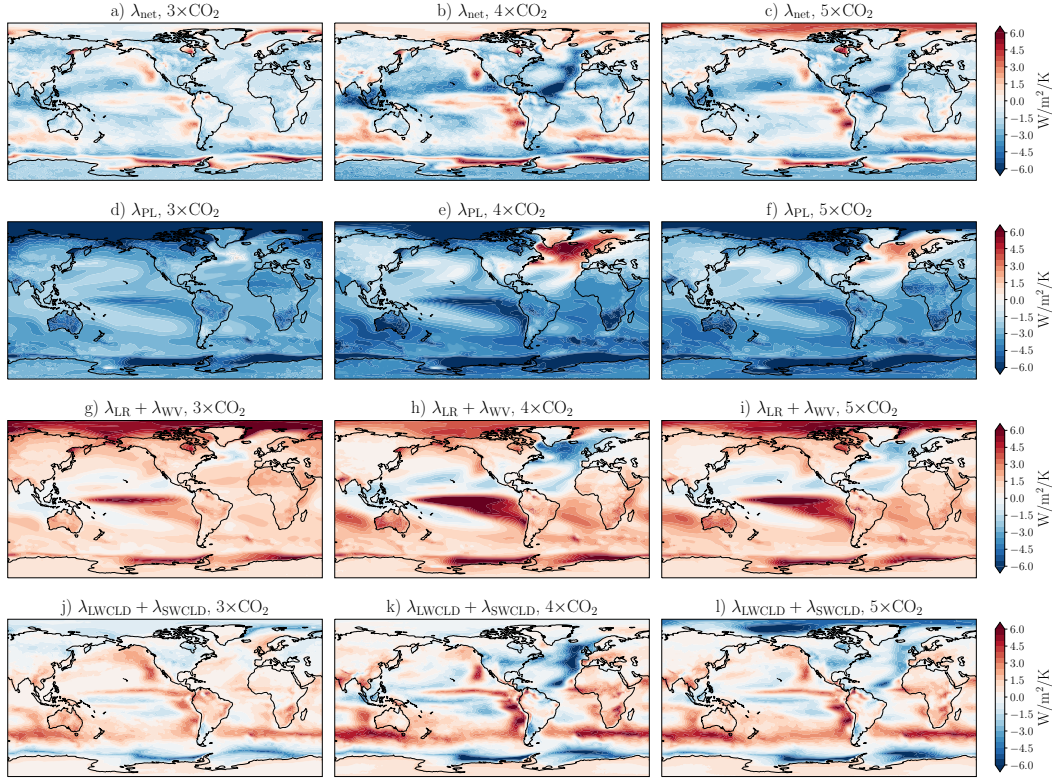
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While the stronger negative feedbacks appear to be located mainly in the North Atlantic, it is unclear whether they are driven by the local North Atlantic SST changes, or by remote SST impacts from other basins. In Fig. 3, we show the normalized SST patterns (over the full 150 years of the simulations) from  $3\times$ ,  $4\times$ , and  $5\times\text{CO}_2$  simulations (panels a-c). We find that anomalous SST cooling primarily occurs in the North Atlantic:  $4\times\text{CO}_2$  produces a strong cooling pattern in the North Atlantic, largely resembling the pattern of the North Atlantic Warming Hole (NAWH) (Chemke et al., 2020). However, this North Atlantic relative-cooling pattern does not emerge at  $3\times\text{CO}_2$  (panel a) and is much weaker at  $5\times\text{CO}_2$  (panel c). Concurrently, we find that local feedbacks exhibit patterns that closely match the SST patterns (Fig. 3f,g, and h). Most of the strengthening of negative feedback that would result in lower EffCS is found at  $4\times$  relative to  $3\times\text{CO}_2$  (Fig. 3i), and it occurs in the North Atlantic, corresponding with the local cooling pattern (Fig. 3d); while most of the weakening of feedbacks at  $5\times$  relative to  $4\times\text{CO}_2$  (higher EffCS, Fig. 3j) which also occurs in the North Atlantic, corresponds with the local warming pattern (Fig. 3e). These results suggest that the non-monotonic response of the feedbacks found in our simulations (Fig. 2a) is predominately from feedback changes in the North Atlantic, associated with North Atlantic local SST changes. We note that significant feedback changes also occur in the tropical Pacific (Fig. 3i,j), particularly the tropical Eastern Pacific, but these feedback changes are in the opposite sign to the global-mean feedback changes, and thus cannot account for the total feedback response we showed in Fig. 2a. While some other regions may contribute to the negative feedback change (e.g., the tropical Western Pacific and the Southern Ocean), we find that the North Atlantic local  $\lambda$  (area between 0 to 60N and 80W

208 to 10E) explains up to 2/3 of the total change in the global-mean  $\lambda$  (Fig. S2). This suggests  
 209 that most of the non-monotonicity at  $4\times\text{CO}_2$  is due to the North Atlantic pattern effect.



**Figure 4.** Maps of individual feedbacks calculated from prescribed-SST runs for: a-c) net, d-f) Planck, g-i) lapse rate + water vapor, j-i) net cloud.

210 To further understand the processes causing the  $\lambda$  non-monotonicity, we further de-  
 211 compose the net feedback parameter  $\lambda$  into the individual feedbacks using radiative kernels  
 212 (Pendergrass et al., 2018) (Fig. 4). In the North Atlantic at  $4\times\text{CO}_2$ , the Planck feedback  
 213 (Fig. 4e) is strongly positive as the local cooling reduces outgoing radiation, whereas the  
 214 combined lapse rate and water vapor feedback (Fig. 4h) and the cloud feedback (Fig. 4k)  
 215 contribute negatively. The strong negative feedback at  $4\times\text{CO}_2$  compared to  $3\times\text{CO}_2$  in the  
 216 subtropical North Atlantic is primarily due to the SW cloud feedback (Fig. S2 and Fig. S3h);  
 217 hence, it is one of the key contributors to the  $\lambda$  non-monotonicity at  $4\times\text{CO}_2$ . This is be-  
 218 cause local cooling strengthens lower tropospheric stability (often measured as estimated  
 219 inversion strength (EIS)), which increases low-cloud cover (negative cloud feedback) and  
 220 more-negative lapse rate feedback. This mechanism is consistent with the leading mecha-  
 221 nism found in the tropical Pacific pattern effect (Zhou et al., 2016; Andrews & Webb, 2018;  
 222 Dong et al., 2019), except this pattern effect here is associated with the North Atlantic SST  
 223 changes (Lin et al., 2019), and causes the non-monotonic response in EffCS and  $\lambda$  across  
 224  $\text{CO}_2$  levels in our experiments. We refer the reader to Lin et al. (2019) for more details on  
 225 this mechanism. Additionally, it is important to note that 2/3 of the difference in feedbacks  
 226 between  $4\times\text{CO}_2$  and  $3\times\text{CO}_2$  comes from the North Atlantic and 1/3 from the rest of the  
 227 globe. At  $4\times\text{CO}_2$ , there are strong responses in the individual feedbacks in the tropical Pa-  
 228 cific (see Fig. S3 for albedo and longwave cloud feedbacks). However, the negative Planck  
 229 feedback response in the tropical Pacific is compensated by the local positive feedback re-  
 230 sponse from lapse rate, water vapor, and clouds (Fig. 4b,e,h,k), which makes the tropical  
 231 Pacific less pronounced in the  $\lambda$  non-monotonic changes.



232 Having shown that feedback changes primarily come from the North Atlantic associated  
 233 with local SST cooling, we finally return to the key pattern of North Atlantic SST cooling  
 234 found in our simulations, the North Atlantic warming hole (NAWH). In the literature, the  
 235 appearance of the NAWH has been attributed to the slowdown in the AMOC and linked to  
 236 an atmospheric response (Rahmstorf et al., 2015; Sévellec et al., 2017; Caesar et al., 2018;  
 237 Latif et al., 2022). Our previous work (Fig. S3 in (Mitevski et al., 2021)) found the North  
 238 Atlantic cooling (NAWH) in our experiments is primarily due to AMOC collapse. The  
 239 AMOC collapses at  $4\times\text{CO}_2$  in our GCM, and at all other higher  $\text{CO}_2$  forcings. At higher  
 240  $\text{CO}_2$  forcings ( $5\times\text{CO}_2$  and above), the AMOC collapse no longer produces anomalous North  
 241 Atlantic cooling compared to the previous level of  $\text{CO}_2$  forcing (e.g.,  $4\times\text{CO}_2$ ) because the  
 242 AMOC collapse-induced SST cooling is further overwhelmed by the surrounding warming.  
 243 Hence, the cooling over the NAWH is less pronounced at higher  $\text{CO}_2$  forcings (Fig. 3c) and  
 244 has a smaller impact on the feedbacks (Fig. 3h). The collapse of the AMOC under  $\text{CO}_2$   
 245 forcing has been widely reported in climate models, including the GISS-E2.1-G model in this  
 246 study (occurring at  $3\times\text{CO}_2$  & higher) and many other CMIP5 and CMIP6 models (Fig. S3  
 247 in (Mitevski et al., 2021)).

## 248 4 Discussion and Conclusion

249 In a series of  $n\times\text{CO}_2$  ( $n = 2, 3, 4, 5, 6, 7, 8$ ) experiments, we find a non-monotonic re-  
 250 sponse in the effective climate sensitivity (EffCS) to  $\text{CO}_2$  forcing using two state-of-the-art,  
 251 coupled climate models. EffCS becomes anomalously low at an intermediate level of  $\text{CO}_2$   
 252 ( $4\times\text{CO}_2$  in CESM1-LE and  $3\times\text{CO}_2$  in GISS-E2.1-G) but increases at higher  $\text{CO}_2$  levels.  
 253 This EffCS non-monotonicity is primarily linked to changes in radiative feedback  $\lambda$  due to  
 254 North Atlantic cooling;  $\lambda$  becomes anomalously negative when cooling emerges in the North  
 255 Atlantic and forms a North Atlantic Warming Hole (NAWH).

256 The dependence of  $\lambda$  on sea-surface temperature (SST) patterns has been widely stud-  
 257 ied, with a focus on the time-evolution of those patterns (Andrews et al., 2015; Zhou et  
 258 al., 2016; Dong et al., 2019; Andrews et al., 2022; Sherwood et al., 2020). For example,  
 259 estimates of EffCS from the observed historical energy budget constraints are lower than  
 260 those from long-term warming under  $\text{CO}_2$  quadrupling, primarily owing to changes in the  
 261 tropical Pacific SST patterns (Zhou et al., 2016; Dong et al., 2019; Andrews et al., 2018,  
 262 2022; Gregory et al., 2020). This “pattern effect” has been studied with a Green’s func-  
 263 tion approach (Dong et al., 2019; Zhou et al., 2017; Zhang et al., 2023), which shows that  
 264 the global feedback has a predominant dependence on tropical western Pacific SSTs and  
 265 is less sensitive to the North Atlantic SSTs. This tropical Pacific SST pattern effect has  
 266 been found to be a leading mechanism for the time evolution of EffCS estimates. However,  
 267 our study proposes a new North Atlantic pattern effect that accounts for changes in EffCS  
 268 and feedbacks across different  $\text{CO}_2$  forcing levels. This North Atlantic pattern effect shows  
 269 that SST cooling in the North Atlantic due to the formation of NAWH causes  $\lambda$  to become  
 270 more negative and, therefore, lower EffCS. We note that the North Atlantic pattern effect  
 271 operates on the dimension of increasing  $\text{CO}_2$  forcing, instead of on the dimension of time  
 272 evolution reported in previous studies (Dong et al., 2019; Andrews & Webb, 2018; Andrews  
 273 et al., 2022; Zhou et al., 2016).

274 The NAWH has been proposed to arise from the reduction of surface meridional ocean  
 275 heat transport (Chemke et al., 2020) or AMOC slowdown that reduces transient warming  
 276 due to increased ocean heat uptake (Caesar et al., 2020; Palter, 2015; M. A. Rugestein et  
 277 al., 2013; Trossman et al., 2016; Winton et al., 2013). In our study, we find that the NAWH  
 278 can further reduce EffCS and transient warming by causing more negative feedback (more  
 279 efficient radiative damping at the top of the atmosphere). The fact that the NAWH has  
 280 been observed in the historical period and is projected to persist in future scenarios with  
 281 increasing GHG (Menary & Wood, 2018; Chemke et al., 2020; Keil et al., 2020; Gervais et  
 282 al., 2018; Liu et al., 2020; Ren & Liu, 2021) suggests a considerable damping effect on future  
 283 global warming from the North Atlantic, which may have been overlooked in the literature.

284 We further analyzed two subsets of CMIP6 models with and without NAWH in the  
285 abrupt-4×CO<sub>2</sub> runs (Fig. S4). Models with NAWH in the abrupt-4×CO<sub>2</sub> scenario also  
286 show more surface cooling in the North Atlantic in transient 21st-century simulations (under  
287 both SSP5-8.5 and SSP2-4.5 scenarios) than models without NAWH. This suggests that  
288 uncertainty in the projected long-term North Atlantic SST patterns in response to abrupt  
289 CO<sub>2</sub> forcing also persists in transient projections. Thus, understanding North Atlantic SST  
290 changes is crucial for constraining global climate change at both transient and equilibrium  
291 timescales.

292 One caveat to our findings is that the AMOC collapse in our models occurs at 3× and  
293 4×CO<sub>2</sub>, which are relatively low CO<sub>2</sub> values, where the collapse can induce a substantial  
294 cooling in the North Atlantic. When the AMOC collapses at a low CO<sub>2</sub> value, such as 2×  
295 or 3×CO<sub>2</sub>, the North Atlantic cooling is strong, leading to a considerable non-monotonicity  
296 in EffCS. However, if the AMOC collapses at a higher CO<sub>2</sub> value, such as 5×CO<sub>2</sub>, then the  
297 overwhelming CO<sub>2</sub> warming from the surrounding areas results in a weaker North Atlantic  
298 SST cooling pattern. In this case, the EffCS non-monotonicity would be smaller than the one  
299 reported in this study. Hence our results suggest that future changes in AMOC and NAWH  
300 may add additional uncertainty to EffCS and transient 21st century warming projections.  
301 Finally, we recognize that our results are obtained using only two GCMs. It would be  
302 important to repeat the same exercise with a broad range of models to test the robustness  
303 of our results.

304 The fact that EffCS is nonlinear and even non-monotonic with respect to CO<sub>2</sub> lev-  
305 els complicates equilibrium climate sensitivity constraints using models, observations, the  
306 paleoclimate record, and process-based understanding. While the non-constant  $\lambda$  across  
307 different CO<sub>2</sub> levels has been mainly attributed to feedback temperature dependence within  
308 models (Bloch-Johnson et al., 2021; Meraner et al., 2013; Mauritsen et al., 2019; Sherwood  
309 et al., 2020; Zhu & Poulsen, 2020) and paleoclimate records (Anagnostou et al., 2016, 2020;  
310 Farnsworth et al., 2019; Friedrich et al., 2016; Shaffer et al., 2016; Zhu et al., 2019), we  
311 here have shown that the SST pattern also plays a role. Our study adds additional evidence  
312 of EffCS state dependence and pattern effects, which need to be further examined in other  
313 lines of evidence.

## 314 5 Open Research Section

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# Supporting Information for “Non-monotonic feedback dependence on CO<sub>2</sub> due to a North Atlantic pattern effect”

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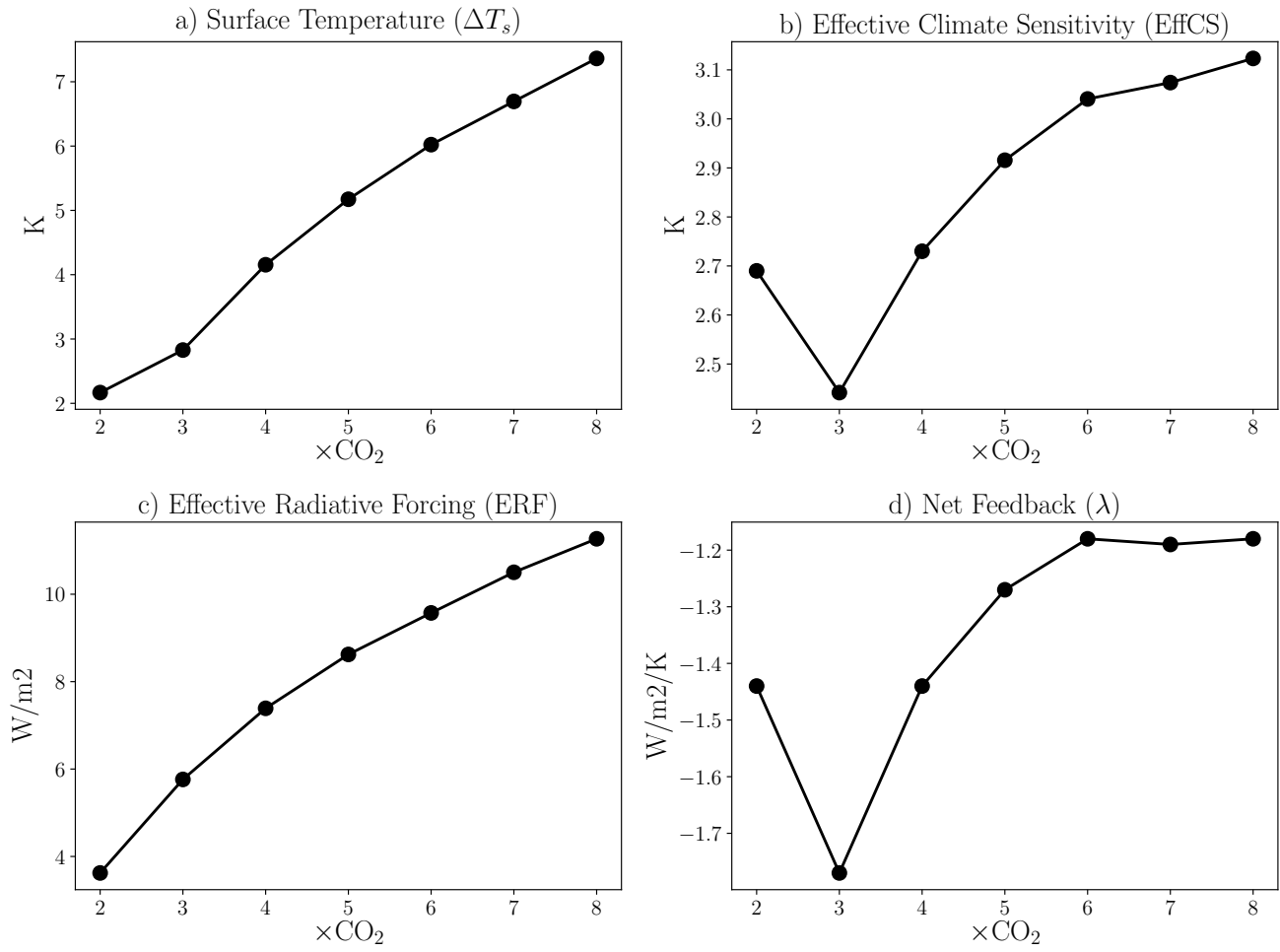
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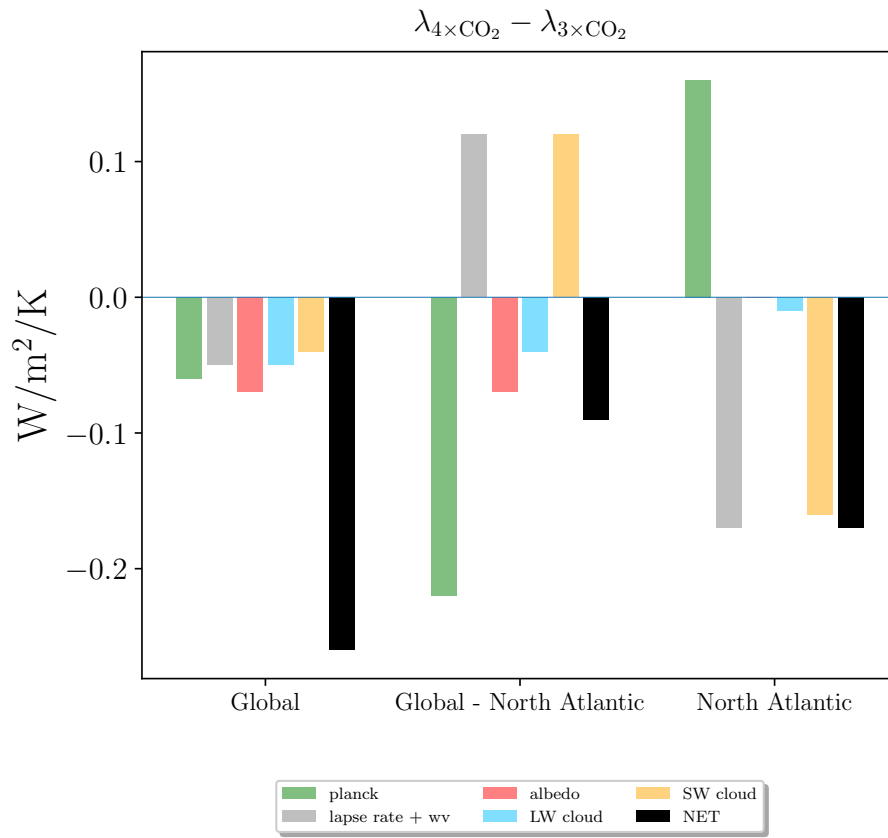
## Contents of this file

Figures S1-S4.

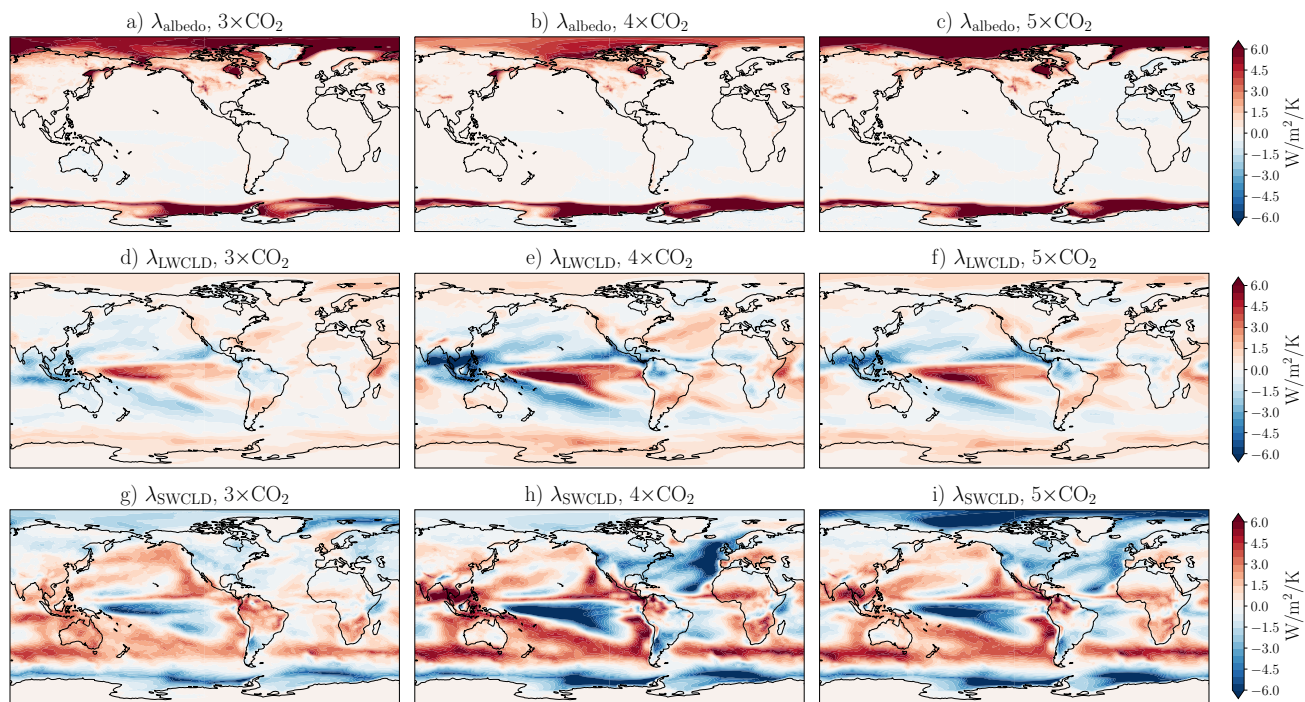


**Figure S1.** Same as Figure 1 in main text but for the GISS-E2.1-G model.

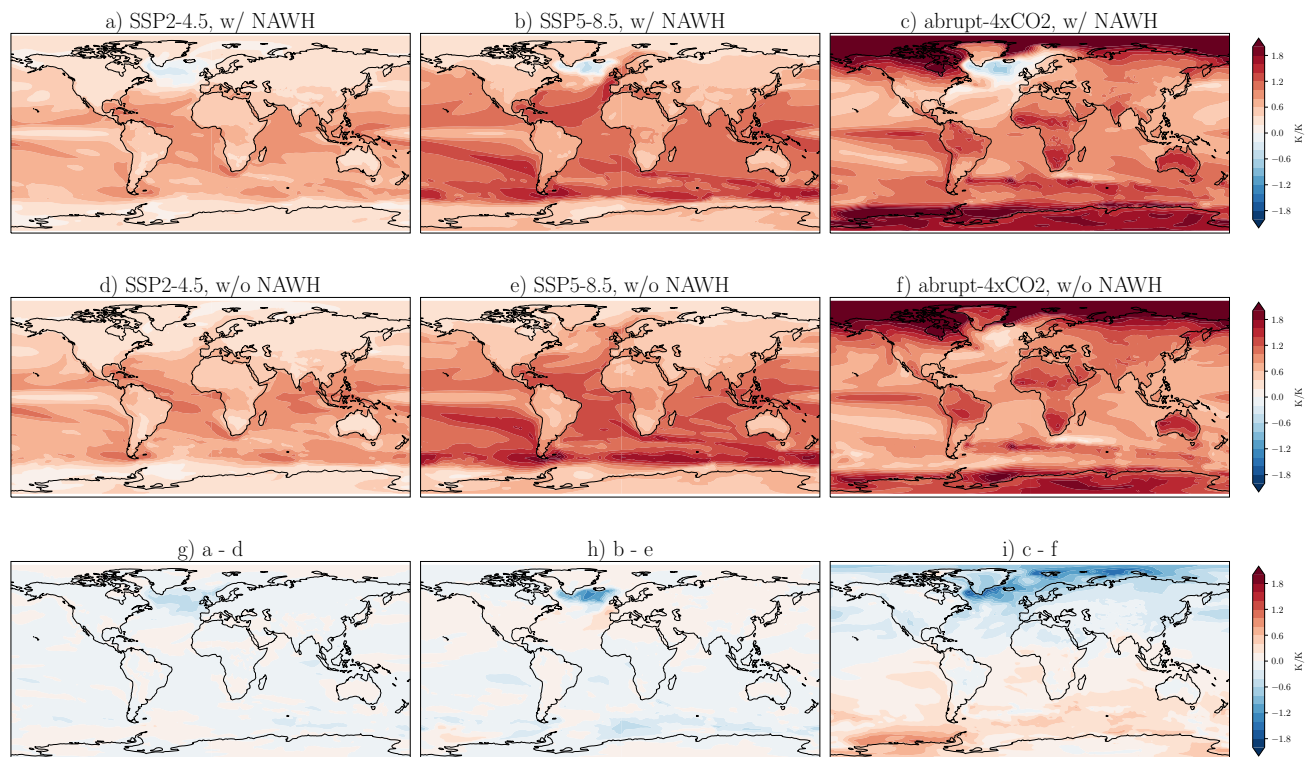




**Figure S2.** Individual feedback difference between  $4\times\text{CO}_2$  and  $3\times\text{CO}_2$  for Global, Global - North Atlantic, and the North Atlantic region. The North Atlantic is defined as a box between 0 to 60N and 80W to 10E.



**Figure S3.** Maps of individual feedbacks calculated from prescribed-SST runs for: a-c) albedo, d-f) longwave cloud, and g-i) shortwave cloud.



**Figure S4.** Maps of surface temperature patterns from two CMIP6 models composites with (a,b,c) and without (d,e,f) North Atlantic Warming Hole (NAWH), defined as cooling in the North Atlantic, and the difference (g,h,i). Composites are shown for SSP2-4.5 (a,d,g), SSP5-8.5 (b,e,h), and abrupt-4xCO<sub>2</sub> scenario (c,f,i). The models without NAWH are ACCESS-CM2, AWI-CM-1-1-MR, CAMS-CSM1-0, CMCC-CM2-SR5, CanESM5, INM-CM4-8, IPSL-CM6A-LR, MIROC6, MPI-ESM1-2-HR, MPI-ESM1-2-LR. Models with NAWH are BCC-CSM2-MR, CESM2-WACCM, FGOALS-g3, GFDL-ESM4, IITM-ESM, KACE-1-0-G, MRI-ESM2-0, NorESM2-MM, TaiESM1. The surface temperature patterns are calculated as local surface temperature changes regressed to global surface temperature response for years 2015 to 2100 for the SSP scenarios and the first 150 years of the abrupt-4xCO<sub>2</sub> runs, and then averaged across models.